

POWERUNIT

EM 2 Notes and Records

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Time Varying Fields

Lecture 1

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$$Vemf = - \frac{d\Delta}{dt}, \Delta = \text{Flux Linkage} = \text{Number of Flux turns}$$

$$= -N \frac{d\phi}{dt}, \phi = \int_{\text{Area}} B \cdot dS$$

$\mu = \#$ of dipoles

⊛ Ferromagnetic Material [Iron & Ferrite]

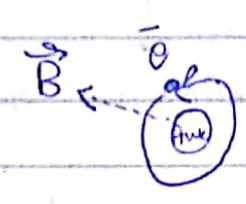
⊛ $\mu \uparrow \uparrow$ Iron $\approx 1500 - 3000$

We use Iron to shield the magnetic field from the External field.

⊛ $\mu \uparrow \uparrow \uparrow$ Ferrite, special characteristics: It's not a conducting material unless you suppress it on itself.

⊛ What's a Dipole? an electron (نورون) & when it does that we get a \vec{B}

$$B = \mu_0 H \quad \text{so} \quad \phi = \int \mu H \cdot dS$$



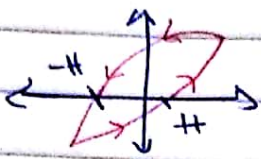
μ is permeability = $\frac{B}{H}$

↳ Forces against the source will induce the μ of -ve particles. μ is the number of these particles.

⊛ applying External magnetic field will rearrange and align to the External field applied [Align the internal fields]

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Each Dipole will generate internal field to support the External one.



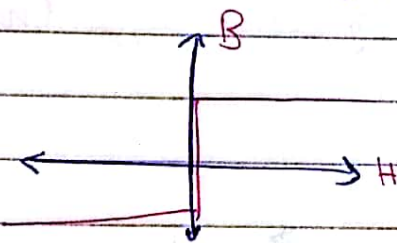
if $\mu \uparrow$ then the # of dipoles will increase.

In order to form the loop from $[-H]$ to $[+H]$ we need $T \equiv$ period because once you remove the applied field, the material will try to go back to its original state \therefore that takes time.

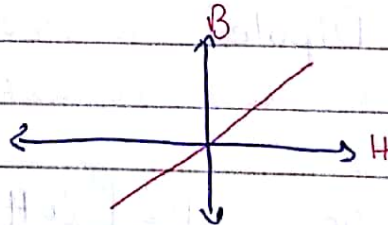
Losses of the Hysteresis are function of frequency.
 $\text{Heat} \propto f^2$

also, μ is a function of frequency.

Ferrite reaches saturation very fast that's why we use air gaps to slow down the saturation.



with no airgaps



with airgaps.

Ferrite has less losses cause there's a free motion but no constricting paths.

Ferrite's μ is very BIG!

Small area in the loop cause small losses
Used in transformers \rightarrow (just like in i(R.S))

⊗ Iron is used to shield magnetic lines

⊗ We tent Iron using \vec{P} or \vec{M}

⊗ $\mu = \sqrt{\epsilon_0 \mu_0} \chi_0$

⊗ $[H]$ as AC from $[-H, H]$

⊗ The loop is a Representation $\left[\begin{matrix} dI/dt \\ \omega I \end{matrix} \right] - [i \omega \mu_0 I]$

↓
equals Losses

⊗ Thermal Energy is Another Representation for "work".

POWER UNIT Lecture 2

$$V_{emf} = -N \frac{d \int B \cdot ds}{dt} = -N \frac{d \int \mu \cdot H \cdot ds}{dt}$$

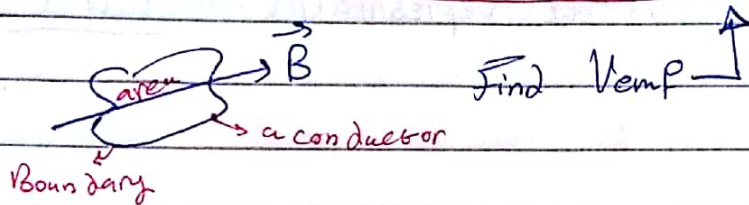
We can interchange between integration and differentiation because they are linear operations and quantities.

$$V_{emf} = -N \int \frac{d(\mu H \cdot ds)}{dt} \quad [\text{Maxwell's equation}]$$

↳ $\int E \cdot dl$ so,

$$V_{emf} = -N \int \frac{d(\mu H \cdot ds)}{dt} = -N \mu \int H \cdot w \cdot \cos(\omega t) ds$$

$$= -N \mu S H \cdot w \cdot \cos(\omega t) = V_{emf} \quad \text{memorize}$$



♡ To limit the boundary and to force charges to move in it. Short circuit the ferritic conductor of the coil.

♡ Electric Field due to Flux linkages

$$\Rightarrow H = H_0 \sin(\omega t) \Rightarrow \text{H level current.}$$

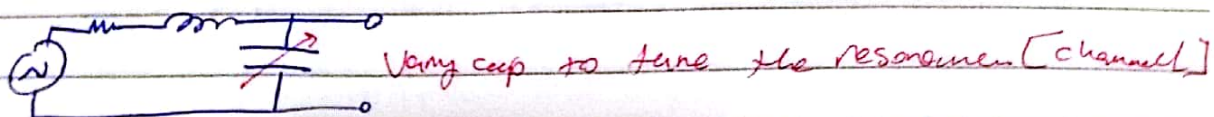
♡ The simplest oscillator is an electron, and around it!

Hard Equations example eg $H_0(x, y, z) = 0$ space

$$= \sinh^{-1} \left(\frac{x}{y^2 + z^2} \right) \quad \text{X-axis to}$$

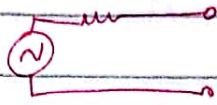
♡ $N \neq$ Must be High in order to receive waves.

♡ The receiver is a band pass filter. [resonance: $\frac{1}{2\pi\sqrt{LC}}$]



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We usually work on resonance frequency so the equivalent circuit will become with resistance only.



Note: Expected Exam Question to find the operating frequency of a microwave with a certain specifications.

⊗ $\lambda = \frac{c}{f}$ mass \rightarrow $\lambda \rightarrow$ $f \rightarrow$

⊗ Higher speed \rightarrow $\lambda \rightarrow$ $f \rightarrow$ current will flow \uparrow
 \rightarrow Eddy current increase \rightarrow more power losses \uparrow

⊗ Small radius \rightarrow Higher frequency.

Micro wave oven operation.

⊕ magnetron [two magnets with high magnetic fields]

⊕ Tripler \Rightarrow Diode + capacitor

⊕ electric field is applied

⊕ Filament [heater] \rightarrow \rightarrow

⊕ but the magnetic field that is applied is very high that keeps them spinning.

⊕ because of the mechanical Design they will spin in Groups (coherent)

⊕ we pick them up using a conductor / usually the voltage is \rightarrow

⊕ Then that voltage is drive to the antenna which will be about 2.45 GHz \rightarrow all the way to the microwave oven room.

⊕ Using Faraday's Cage to shield the room.

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Time Varying current is related to magnetic field that affects induction which is simply Maxwell's equation.

Volume current density $\vec{J} = \sigma \vec{E}$ → Electric field
Conductivity

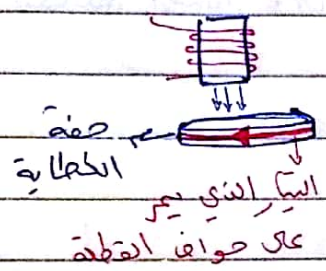
Applications on Faraday's Law

Cap Sealing

Center of Gravity Lab

The intensity in the coil

$J \equiv A/m^2$



outer area is the largest

So, $\int \vec{E} \cdot d\vec{L}$ will increase

and because the foil resistance is

very low the power $P = I^2 R$

will increase as well. As heat

will melt down the disk → (المغناطيسية)

Area ↑, Voltage drop ↑

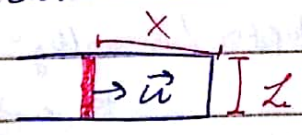
$\mu \uparrow$, High frequency

all of these cause to

intensity the magnetic fields

Electric braking

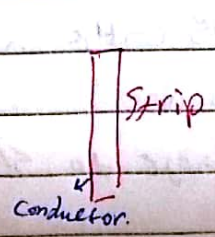
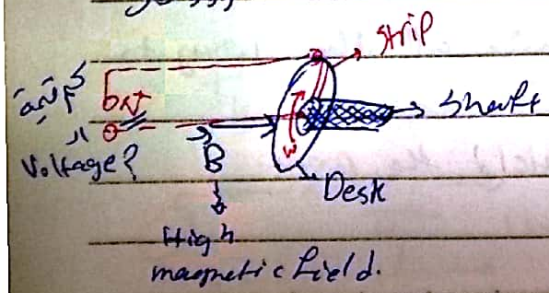
recall



$\mathcal{E} = X \cdot L$, $V = -B w L$

$w = w \cdot r$, distance from center

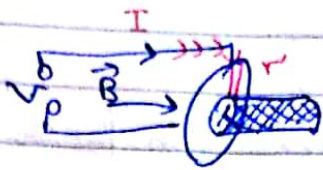
So...



$\int du = \int_0^R r B w dr$

$\int du = \int_0^R B w r dr \Rightarrow V = \frac{w B R^2}{2}$

if we were to find the current??
 $V = \frac{\omega B r^2}{2}$ $I = \frac{V}{R}$ (short circuit)



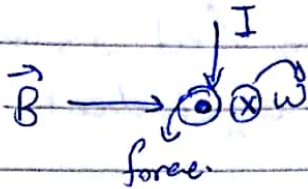
↓ I ... so force ↓
 , (I) e

♥ magnetic field induced is always in the opposite direction of the source that caused it.

♥ The current will radially move in the disk.

♥ if we have magnetic field + current = Force

Direction = $I \times B$



So, the Braking will happen.

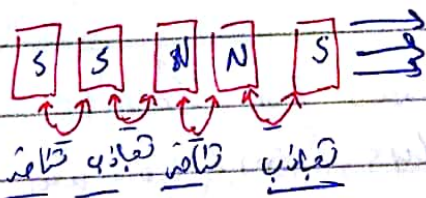
if R Load ↓, Braking Force ↑↑

⊗ Examples

- 1] wheel: Fly wheel [electric braking system]
- 2] Hybrid cars: (can generate energy)

3] to convert the ^[motor] load into a generator, we must apply force with higher speed than the field
 we usually use AC Generators (synchronous generator)

3] Linear motor.



So, it will move!

to change speed, change frequency
 ↑
 Related.

eg Lift the train above ground.

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Lecture 4

☞ ϵ_0 dipole reorientation / # of = zero

☞ σ for insulation

☞ μ_r for Ferrite is largest

☞ $\mu_r \uparrow \uparrow$ (linearly) $\rightarrow v \uparrow \uparrow$ (linearly)

Linearly [Like a multiplication factor for both side source & affected]

☞ $\mu_r \Rightarrow$ Used to intensify the magnetic field lines

☞ μ is nonlinear (problem) force for magnetic field \uparrow electric field \downarrow [Flux]

☞ μ is linear when the magnetic field is small.

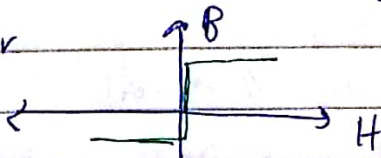
☞ Real power (in thermal form) is equivalent to losses

"Resistance" but depends on frequency. not a normal Resistance (ohmic losses)

☞ When free electrons you increase conductivity # of free electrons \uparrow \rightarrow conductivity \uparrow

☞ Breaking down due to high magnetic external field or force.

☞ Ferrite is used in logic gates such as Schmitt trigger



Hysteresis for Ferrite

in this form \times so we add air gaps in order to use it in transformers.

we can Amplify the signal / voltage when we increase the turns Ratio.

It can also be used as a differentiator.

σ = electron motion Representation.

σ goes to Zero $\rightarrow R$ increases \rightarrow insulator.

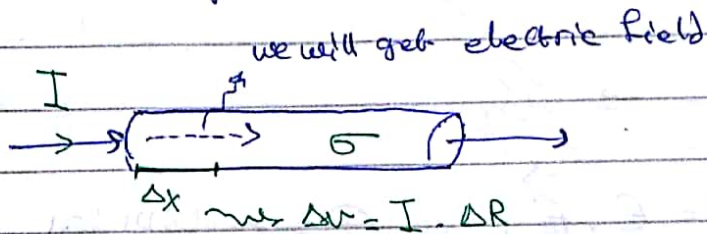
σ goes to ∞ $\rightarrow R$ decreases \rightarrow conductor.

Any space or material will have $[E, \sigma, \mu]$

\downarrow
electrons
capability
to store
energy.

ϵ_0 can be obtained statically
by applying different potentials

$$E = \frac{\Delta V}{\Delta x}$$



$$J = \sigma E$$

Electric field ϵ of conductor or semi-conductor μ is ϵ_0 conductor.

Which is the basic idea, conductor is like dipole for antenna.

From
If we apply External field $[+q \rightarrow -q]$
force will be generated to move \vec{e} .

This is called polarization.

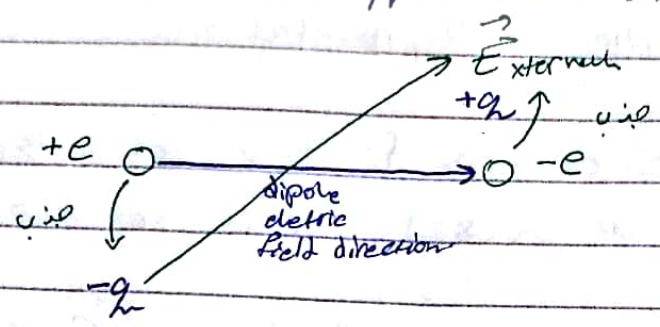
When applying force distance between \vec{e} & the charge change \rightarrow balance will change \rightarrow polarization happen!

[e.g. liquid crystal]

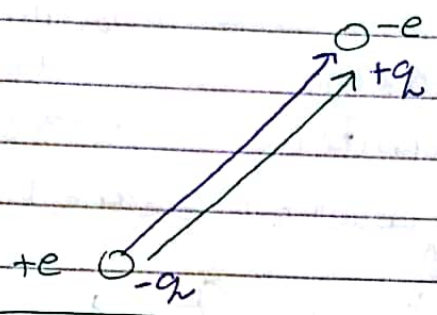
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Image = Reflection of charges.

Sensors are an application on electromagnetism



Reorientation



$$D = \epsilon \cdot E$$

Charge displacement

$\epsilon = \epsilon_0 \epsilon_r$ ϵ_r is the relative permittivity of the material

Dielectric: $\epsilon_r > 1$, $\epsilon_r \uparrow \Rightarrow$ electric field $\downarrow \downarrow$
 when dipoles move \Rightarrow external field.

Real Energy: when we apply work and bonds are broken!

$$\epsilon_r = \epsilon_r' + j \epsilon_r''$$

real power loss

mainly as a Resistance
 can be seen in Microwave ovens
 a characteristic in water when
 we apply STRONG external
 field.

$$\mu_r = \mu_r' + j \mu_r''$$

Expressed in eddy currents
 [rare]

Wave Equation

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$\oint E \cdot dl = -N \omega \mu H_0 \cos(\omega t)$ ← single harmonic case.
 ↑ single point in Fourier domain.

orthogonality: ① inner / dot product = zero

② cross product direction (موجة دائرية في المستوى)

Maxwell's Equation again.

$$\nabla \times E = -\frac{dB}{dt}$$

∇ = direction (اتجاه) \rightarrow ! issue

→ For directive quantities.

e.g: $E = E_x \vec{a}_x = E_x(x, y, z, t) \vec{a}_x$ direction.

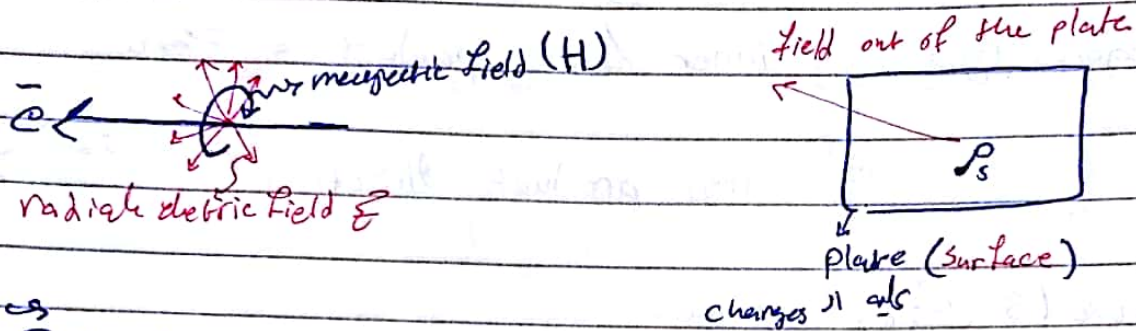
Value ←

⊗ Value could be a function of anything (in e.g, space x, y, z, t)

⊗ If we hold these parameters (القيم) E becomes a constant (Real Number).

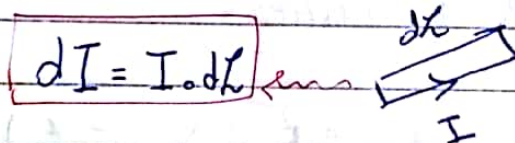
☞ Heat \equiv Infrared radiation.

☞ For the same source (in space) electrical field and magnetic field are orthogonal.



☞ $E \propto \frac{1}{R^2} \propto q \propto \frac{1}{\epsilon_0}$

☞ charge (current element) \Rightarrow field generation



☞ Simplest Maxwell's Equations ☞

$$\nabla \times E = -\frac{dB}{dt}$$

$$\nabla \times H = J + \frac{dD}{dt} \quad , \quad D = \epsilon E$$

also, $J = \sigma \cdot E$

$$\nabla \times H = J + \epsilon \frac{dE}{dt}$$

$$\nabla \times H = \sigma \cdot E + \epsilon \frac{dE}{dt}$$

☞ we propagate the electromagnetic waves in space.

For free space $\nabla \times H = \epsilon \frac{dE}{dt}$

☞ $\nabla \times$ \Rightarrow it's a derivative for orthogonal stuff, also some sort of summation.

Q3 We Express cycling arguments in time and space

$E = E_0 \sin(\omega t - \beta r)$, How the electric field should look like

direction of propagation = \vec{a}_r (unit vector)

$H = H_0 \sin(\omega t - \beta r)$, How the magnetic field should look like.

$\omega \equiv$ frequency in time = $\frac{2\pi}{T}$

$\beta \equiv$ frequency in space = $\frac{2\pi}{\lambda}$, λ is the period in space or wave length

Direction of propagation

$\vec{a}_E \times \vec{a}_H = \vec{a}_r$

e.g. if $E \rightsquigarrow \vec{a}_x$ & $H \rightsquigarrow \vec{a}_y$

$E = E_0 \sin(\omega t + \beta z) \vec{a}_x$

$H = H_0 \sin(\omega t + \beta z) \vec{a}_y$

Electric field \times magnetic field
cross



Gives you the direction of power carried on the electro magnetic wave.

So $\vec{a}_x \times \vec{a}_y = \vec{a}_z$

$\gamma = z$ when we rewrite the answer in the argument we take the vector out

\rightsquigarrow because we send (power density) and Energy through transmission!

$H = H_0 \sin(\omega t + \beta z)$
not \vec{a}_z

Plane wave in Free space

$$\left. \begin{aligned} \star E &= E_0 \sin(\omega t + \beta r) \\ H &= H_0 \sin(\omega t + \beta r) \end{aligned} \right\} = \text{Star} \star$$

origin of wave equation:

$$\nabla^2 E - \gamma E = 0$$

del operator.

Relation between E & H

$$E = \frac{\omega}{\beta} H, \quad \frac{\omega}{\beta} (\text{air}) \text{ median characteristics impedance (could be complex)}$$

[proportional to voltage signal]

[proportional to current signals]

Direction of propagation is a unit vector (γ)

$$E \times H = \gamma$$

All waves can be approximated to (plane wave)

How to achieve plane wave? when the electric field is a straight line on a certain axis & the magnetic field is on a different orthogonal axis and they are both from the same source.

From star \star you can tell that H_0 & E_0 are much larger than the source.

propagation direction tells us where the source's at.

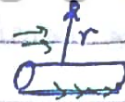
While the wave propagates \rightarrow it will lose power \rightarrow lose H
 \rightarrow lose E

Boundary condition

to create electric field in a wire you need electric field on the outer space (air) surrounding this wire.

to create electric field outside the wire you need electric field inside it.

Boundary condition: The tangential component of the electric field is always continuous.



power density or Energy $\equiv \text{WATT/m}^2$

as a function of distance E_0 will decrease/reduced by $\frac{1}{r^2}$.

$\frac{1}{Z} = \sqrt{\frac{\mu}{\epsilon}}$ (a real quantity)

medium characteristics impedance for free space.

From Bohr's Equation

$$u = \frac{\omega}{\beta} = \frac{\frac{2\pi}{T}}{\frac{2\pi}{\lambda}} = \frac{\lambda}{T} = f \cdot \lambda \quad \text{also} = \frac{1}{\sqrt{\mu \cdot \epsilon}} = \frac{c_{\text{speed of light}}}{\sqrt{\mu_r \epsilon_r}}$$

Speed of Propagation

$\rightarrow \mu = \mu_0, \epsilon = \epsilon_0$

then u is maximum.

if we fix the [medium / material / environment] then we fix the speed of propagation.

if we fix the frequency [not a function of time / space] then we can measure λ

if we fix time [by taking more than one point at the same time] then electric field will be sinusoidal

Nyquist criteria: to represent a sinusoid you need at least two [2] points.

Remotely, we can measure the source by the electromagnetic waves

Transmitting Antenna

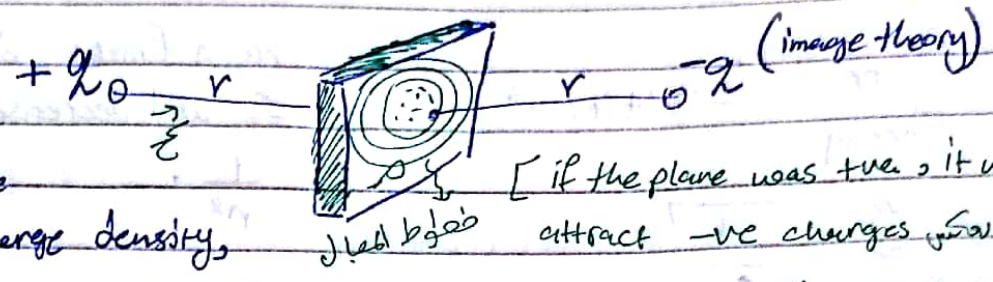
if we pass electric field in a wire this electric field will show up in air surrounding the wire.

Receiving Antenna

and if there's an electric field outside the wire [in air] an electric field will be induced in the wire.

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Q What ever we see is basically $\rho =$ charge density.
 Charge density reflect the source [usually light]



Q we only see the surface charge density, which is \equiv charges ρ_s Surface. ρ_s is [if the plane was true, it will attract -ve charges equally]

Q Anything that Reflect images is of [conductor]
 conductivity ρ \rightarrow ρ \rightarrow ρ \rightarrow ρ
 [good conductor, semi conductor, ...]

σ, ϵ, μ A function of frequency.

Q $\frac{1}{Z} \neq$ Heat losses ρ \rightarrow ρ \rightarrow ρ \rightarrow ρ \rightarrow it's an impedance, not a Resistance.

Q Total Real power can be obtained from EXH 's value indicates total power ρ stored.

Q if $\frac{1}{Z}$ was complex

Q Real part ρ (lossless) will not dissipate any real power

Q Imaginary part ρ (heat losses) will consume power.

VERY IMPORTANT circuits ρ \rightarrow ρ \rightarrow ρ \rightarrow ρ

Q λ (wave length) is measured in (meters).

Q the charge doesn't propagate but its effect does! [electromagnetic wave] ρ \rightarrow this weight [photons].

$$E(x; t) \Rightarrow E_x(t) = E_0 \sin(\omega t - \beta z) \vec{a}_x$$

$$H(y; t) \Rightarrow H_y(t) = H_0 \sin(\omega t - \beta z) \vec{a}_y$$

$E \times H = \gamma$ propagation axis (+ve z-axis)

$$E_0 = \sqrt{\mu} H_0 \text{ (in lossless medium)}, \sqrt{\mu} = \sqrt{\frac{\mu_0}{\epsilon_0}} = 120\pi = 377 \Omega$$

if $r = -z$

$$\rightarrow E = E_0 \sin(\omega t + \beta z) \vec{a}_x$$

$$\rightarrow H = H_0 \sin(\omega t + \beta z) \vec{a}_y$$

Relation between (time) and (space) is speed.

$$\lambda \cdot f = u \text{ (free space where } \mu = \mu_0, \epsilon = \epsilon_0)$$

$$\frac{1}{\sqrt{\mu_0 \epsilon_0}} = c \text{ (speed of light)} = 3 \times 10^8 \text{ (maximum)}$$

Media types

Free space [air, $\mu = \mu_0, \epsilon = \epsilon_0$]

even if there was no material at all, the electromagnetic field will be supported by the interaction between them.

Free space will not be affected by electric field

nor magnetic field.

Theoretically we don't let the intensity of electromagnetic field change the physical characteristics of the material

material's properties such that electromagnetic field will be within normal field

values.

$$E = 10 \sin(2\pi \times 10^9 t - \beta z) \vec{a}_x$$

$$\lambda = \frac{3 \times 10^8}{10^9} = 0.3 \text{ m}$$

$$\beta = \frac{2.0\pi}{0.3}, u = \frac{\omega}{\beta}, \beta = \frac{\omega}{f \cdot \lambda}$$

$$H = \frac{10}{120\pi} \sin(2\pi \times 10^9 t - \beta z) \vec{a}_y$$

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⊕ We also need to find the power density:

$\vec{E} \times \vec{H}$ \rightarrow magnitude only has the power density but the direction points in the propagation axis. Hence it is called a (pointing vector).

$$|P_0| = \frac{1}{2} E_{\max} \cdot H_{\max} \cdot \sin(\theta) \quad (\text{W/m}^2)$$

in between.

in our case \rightarrow plane wave they are orthogonal

$$|P_0| = \frac{1}{2} E_{\max} \cdot H_{\max}$$

⊗ How to generate an electromagnetic wave?

\rightarrow From the boundary condition: [1] conductor

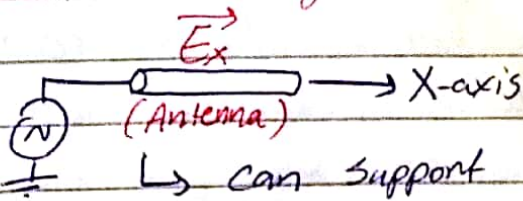
[2] time varying current

⊗ if the current frequency was 1 GHz \rightarrow the wave frequency will also be 1 GHz

will produce electric field in the conductor \rightarrow so electromagnetic wave surrounding that conductor will be established!

⊗ To change the wave frequency, the only way is to change the supply (current)'s frequency.

The function diagram for this case:



\rightarrow can support current & magnetic field

⊗ We can pick an electromagnetic wave from a steady step (sharp @ maximum component) using an antenna that is tuned at a certain impulse.

⊗ Pulse frequency contents are about (∞) A wide spectrum.

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⊗ Antenna size (along with almost all electrical devices) depends on λ (meters) & can be controlled by frequency.

$$\text{Cause } \lambda = \frac{c}{f}$$

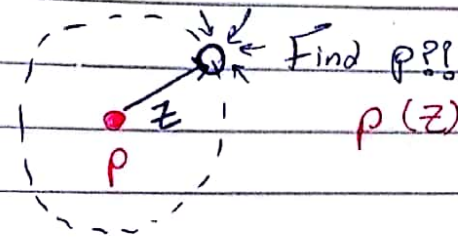
⊗ e.g. For a station @ 801 kHz find $\lambda = 371 \text{ m}$

⊗ tower height $\Rightarrow \frac{\lambda}{4} \approx 91 \text{ m}$

⊗ tower height $\times 2 =$ distance between towers. $= \frac{\lambda}{2} = 182 \text{ m}$

⊗ Mainly, electromagnetic wave power depends on exactly how much you are sending and receiving (circuits)

⊗ calculating power at a certain point ⊗

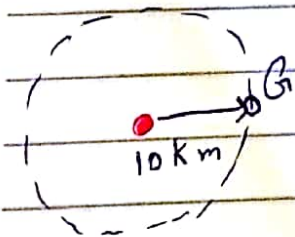


$$P(r) = \frac{\text{Power transmitted}}{\text{distance}} = \frac{P_t}{4\pi r^2} \text{ (W/m}^2\text{)}$$

$\hookrightarrow (r)^2$

which implies $P_D = \frac{1}{2} E_m \cdot H_m$ ←

Ex → if $P_t = 1 \text{ watt}$, find P_D @ G :-



$$P(G) = \frac{1}{(10 \times 10^3)^2 \cdot 4 \cdot \pi} = 795 \text{ p (watt/m}^2\text{)}$$

a very small value!

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Ex →

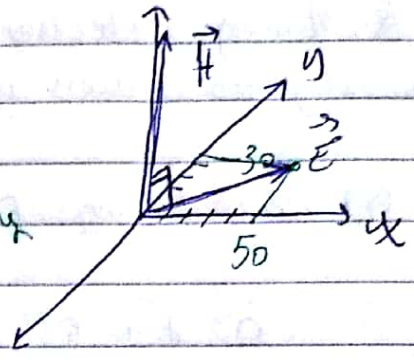
$$E = 50 \sin(\omega t - 30\pi z) \hat{a}_y - 30 \sin(\omega t - 30\pi z) \hat{a}_x$$

Find frequency Draw the curl (Relation) of H?

$$H = \frac{50}{120\pi} \sin(\omega t - 30\pi z) \hat{a}_x - \frac{30}{120\pi} \sin(\omega t - 30\pi z) \hat{a}_y$$

Linear (two dependant component)

H is orthogonal on Electric field.



Ex → give two orthogonal lines for the following in the same plane?

$$a \hat{a}_x + b \hat{a}_y$$

- 1) $-b \hat{a}_x + a \hat{a}_y$ } orthogonal
- 2) $b \hat{a}_x - a \hat{a}_y$ }

in phase, same frequency, speed

Coherently, they vary as a function of time and space.

(the two parts)

⊗ If β was $+30\pi$ in the previous Example or they were not in phase or the direction of propagation was different, then they are not from the same source and they are not components for the same wave.

⊗ if \vec{E} had \uparrow x, y & z components then they are \uparrow the same wave and this is a vector representation.

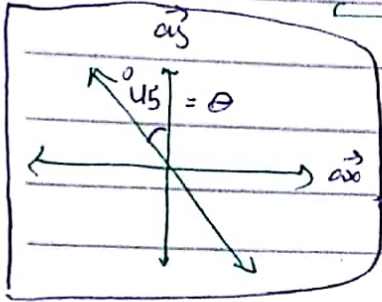
⊗ If the signals/waves were orthogonal then:

- 1] the dot product is equal to zero. } we use this when
 - 2] the cross product is equal to one. } we have unknowns
- we "dot product" them
them = 0, solve in

⊗ Polarizations Represents the trace of the tip of the electric field vector.

Why trace? it changes in space & time. (gives an argument of sine & cosine.)

$$E = E_0 \sin(\omega t - \beta r) \vec{a}_r$$



- ⊕ Because the argument doesn't affect anything
- ⊗ This obtained line will not change if we change the values cause it's an argument.
- ⊗ The electric field will change on this line
- ⊗ called linearly polarized wave جانبياً

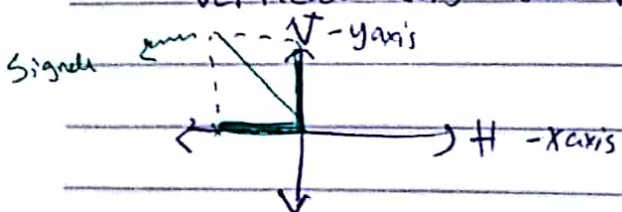
⊗ orientation of receiving or transmitting antenna shape depends on the polarization so the tangential component of electric field outside goes in the conductor so we can receive it.

⊗ Waves in space can be supported if we give them two orthogonal parts (H, V)
Horizontal
Vertical

⊗ if the angle θ between them is 90° , it (Antenna) will not receive anything

⊗ Inner product is the one that receives stuff, that's why the conductor of the Antenna should be orthogonal with the electric field lines.

⊗ Horizontal axis will receive (الجزء الأفقي) .
 Vertical axis will receive (الجزء العمودي) .



⊗ هذا الجزء الأفقي هو الجزء الأفقي
 Inner product (cosine)
 ! two axis signal

POWER UNIT Lecture 11

When $\sigma = 0$ then there is no attenuation due to conduction.
 Also, No drop in E (Electric Field) & No R (Resistance)
 So No heat Losses.

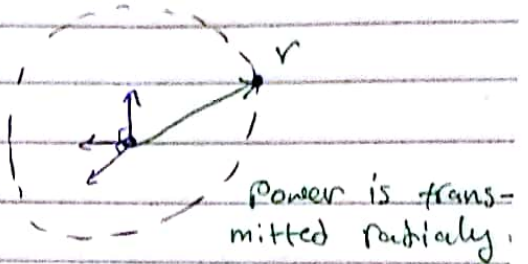
Two types of Losses:

- 1] Due to distance
- 2] Due to media / material

⊗ Losses due to distance :-

$$P_D = \frac{P_t \text{ transmitted}}{4\pi r^2}$$

Power Density



$$P_r = A_{\text{eff}} \cdot P_D$$

received effective Area

⊗ Attenuation due to space.

$$P_r = \frac{P_t \cdot A_{\text{eff}}}{4\pi r^2} \quad \text{called (Link budget) which is the relation between } P_r \text{ \& } P_t$$

⊗ Losses due to material.

$\sigma \neq 0$ (Lossy media)

$$\Rightarrow \nabla^2 E - \gamma E = 0 \quad \text{maxwell's / Laplace Equation}$$

$$\Rightarrow \gamma^2 = j\omega\mu (\sigma + j\omega\epsilon)$$

$$\Rightarrow \gamma = \alpha + j\beta \quad \text{where } \beta \text{ from previous lectures.}$$

$$\Rightarrow \alpha = \omega \cdot \sqrt{\frac{\mu\epsilon}{2} \left[\sqrt{1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2} - 1 \right]}$$

$$\Rightarrow \beta = \omega \cdot \sqrt{\frac{\mu\epsilon}{2} \left[\sqrt{1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2} + 1 \right]}$$

POWER UNIT

3) Steady State Solution

→ $\vec{E} = E_0 \underline{e^{-\alpha z}} \sin(\omega_0 t - \beta z) \hat{a}_z$ direction of propagation.

• Now $e^{-\alpha z} \Rightarrow$ called the attenuation factor.

→ if there was no propagation, energy will decrease which cause more power loss.

⊗ Real power is lost as heat in the media \rightarrow conduction path / current \rightarrow infrared region.

⊗ if $\alpha \uparrow \rightarrow$ losses will $\uparrow \rightarrow$ more heat $\uparrow \uparrow$.

⊗ (H) direction doesn't change

→ $\underline{H} = \frac{E_0}{\eta} e^{-\alpha z} \sin(\omega_0 t - \beta z) \hat{a}_y$

⊗ For the same wave, the sine argument is the same for \vec{H} & \vec{E}

⊗ Now what is η ??

$\eta = \sqrt{\frac{j\omega\mu}{\sigma + j\omega\epsilon}} \Rightarrow |\eta| \angle \theta_\eta \Rightarrow |\eta| \cdot e^{j\theta_\eta}$

take j to the denominator $\rightarrow -j$ & the existing j will be gone.

→ $\frac{j\omega\mu}{-j\sigma + \omega\epsilon}$ so, $\theta_\eta = \frac{1}{2} \tan^{-1} \left(\frac{\sigma}{\omega\epsilon} \right)$

and $\frac{\sigma}{\omega\epsilon}$ is called the loss tangent.

Which is one of the material's characteristics!

POWER UNIT

If $Z = 50 \angle 20^\circ$

then there are losses & we know that cause there is an angle & it's a complex quantity

⊙ if θ_{xy} exist, How can we write H?

$$H = \frac{\epsilon_0}{|Z|} e^{-\alpha z} \sin(\omega t - \beta z + \theta_{xy})$$

 constant phase shift due to media material.

⊙ α Represents the attenuation and it usually appears at low frequencies, air that is close to earth (MF and below) (that's why it's called surface wave propagation attenuation). at High frequencies, α is almost gone.

⊙ The Higher the Frequency the straighten the propagation.

exists ↙ ↑ →	ELF	0.3k - 3k	} Why the [3]?! because it's easire with λ equation $\lambda = \frac{3 \times 10^8}{f}$	
	VLF	3k - 30k		
	LF	30k - 300k		
	MF	300k - 3M		
	HF	3M - 30M		
	VHF	30 - 300M		
	UHF	300M - 3G		
EHF	3G - 90G			

At High Frequency the ionosphere will reflect the wave (called sky wave)

⊙ VHF and above, the wave will escape the ionosphere. [causes random rotation to the propagated waves]

3 Perfect Conductor (μ & ϵ have no effect)

σ goes to ∞

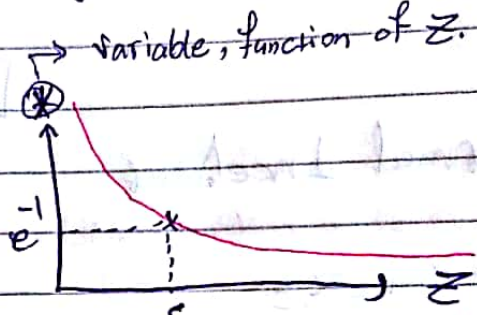
The propagation of electric field in a conductor is zero.

There's difference between the propagation & existence of the electric field.

When we apply current to a conductor, electric field exist to support the electric field that already exist in the conductor, as for the wave propagation of the electric field, it is zero!

$E = E_0 e^{-\alpha z} \sin(\omega t - \beta z) \vec{a}_x$

negative weight can be neglected



$$\alpha = \omega \sqrt{\frac{\mu \epsilon}{2} \left[1 + \left(\frac{\sigma}{\omega \epsilon} \right)^2 \right]}$$

Goes to ∞ can be neglected

$$\therefore \alpha \approx \sqrt{\frac{\mu \sigma \omega}{2}} \rightarrow 2\pi f$$

$$\alpha \approx \sqrt{\pi \mu \sigma f} \quad \text{So,}$$

$$\delta = \frac{1}{\alpha} = \frac{1}{\sqrt{\pi \mu \sigma f}}$$

δ delta = a small value
skin depth

We say, @ δ the E on surface is e^{-1} . So,

$$\delta = \frac{1}{\alpha}$$

Finding New α in conducting materials.

Thin & thick conductors (expected Quiz Question)

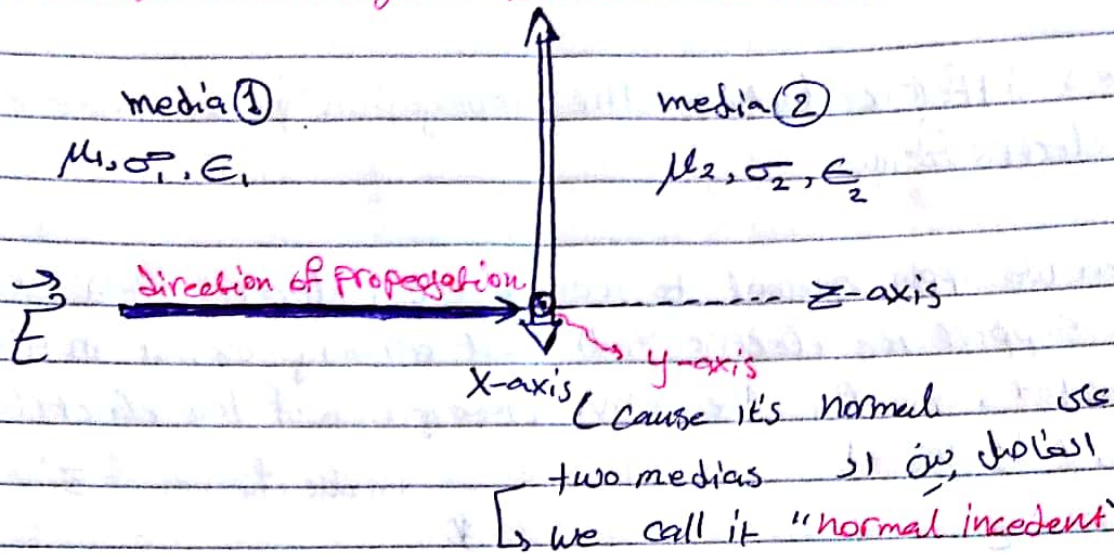
- Thin: when it's comparable with the skin depth \rightarrow *مقارب*
- Thick: when it's about 5-2 times the value of the skin depth \rightarrow *بأضعاف*

POWERPOINT

Applications on thin & thick conductors: [Shielding]

Incidents and Reflection

if we had an electromagnetic wave in a media:



Normal Incident

(The most important case • Applications \Rightarrow Transmission Lines)

$$\vec{E}_i = E_0 e^{i(\omega t - \beta_1 z)} \vec{a}_x \quad \text{in General!}$$

The boundary conditions

\Rightarrow tangential component is the boundary between the two materials.

\Rightarrow the normal component is normal on the boundary.

⊗ difference in Voltage drop causes current

(tangential component) \Rightarrow normal component

The tangential component is Always continuous.

$$E_{t1} \Delta x = E_{t2} \Delta x, \text{ so}$$

tangential

$$E_{t1} = E_{t2}$$

POWER UNIT

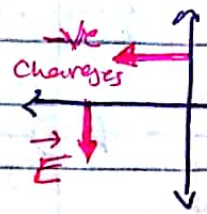
$D_{n2} - D_{n1} = \rho_s \Rightarrow$ if perfect conductor then $[D_{n2}]$ is zero, so ρ_s is negative

$D \equiv$ flux lines (normal)

$\rho_s = -D_{n1}$

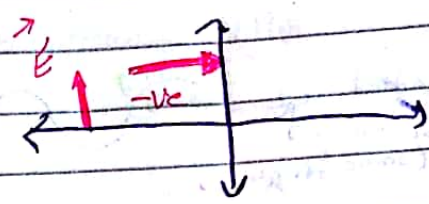
⊗ ρ_s Represents the Reflecting image of the source of energy.

if \vec{E} is direction



negative charges direction is as shown.

if \vec{E} direction changes



⊗ If we had a Dc current that is passing through a perfect conductor then the propagation is (Zero) and this happens in a very short time. Once we Apply External field the material Reconfigure its electrons in a period of time (Relaxation time)

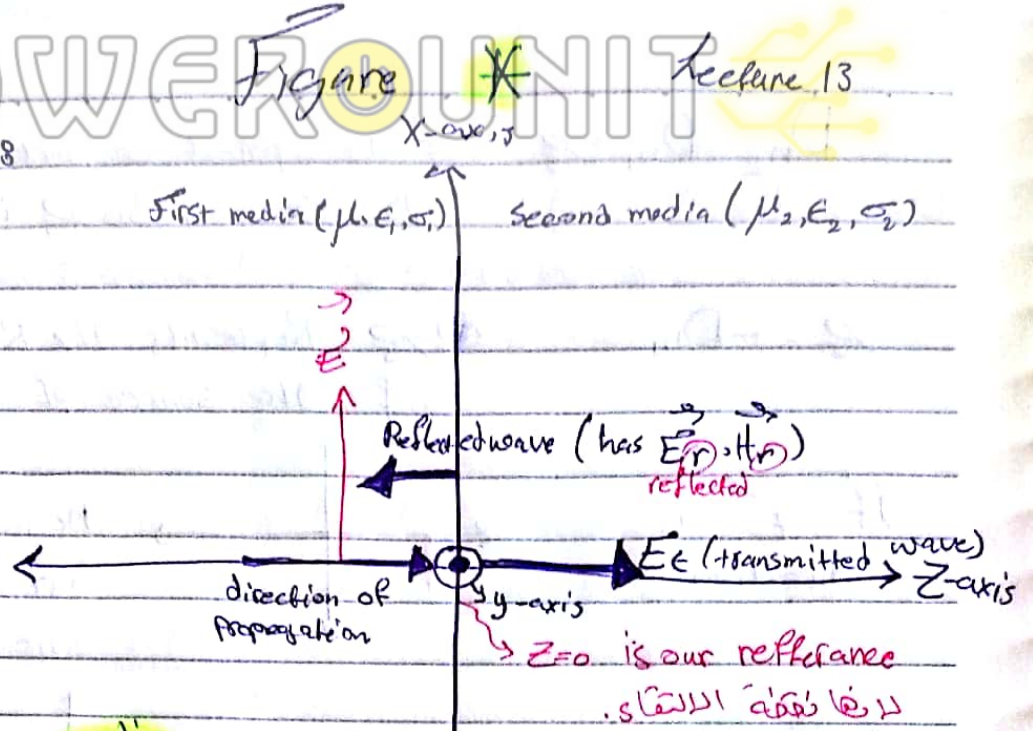
المادة لا تملك وقت الاستجابة في الـ material

When Relaxation time increases, Delay in Reflecting an image (surface charge density ρ_s) increases.

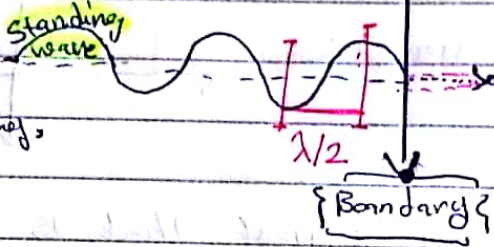
⊗ For conductors, practically we say that there's no normal flux lines, because the relaxation time is very small, they are killed fast. That's why we don't draw them.

POWER UNIT Lecture 13

⊗ Neglecting α 's



The total E added two waves because they are coherent (same frequency).



not a real wave, it is the effect of superposition on the incident and reflected wave in arbitrary point z

First media

- $E_x = E_0 \sin(\omega_0 t - \beta z) \hat{a}_x$
 - $H_y = \frac{E_0}{Z} \sin(\omega_0 t - \beta z) \hat{a}_y$
 - $\beta_1 \neq \beta_2$ (different medias)
- Also, λ is different in different media, so...

• $U_1 \neq U_2$
speed of propagation

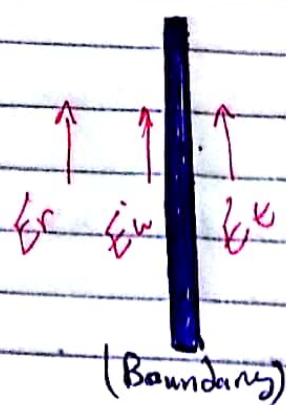
$$U_1 = \frac{1}{\sqrt{\mu_1 \epsilon_1}} = \frac{c}{\sqrt{\mu_r \epsilon_r}}$$

Second media after reflection

- $E_r = E_i * \Gamma$ \rightarrow same direction
- $E_i \equiv$ incident $\left[\begin{matrix} \text{is absolute} \\ \text{Boundary} \end{matrix} \right]$
- $\Gamma = \frac{Z_2 - Z_1}{Z_2 + Z_1}$ \rightarrow could be complex

Δ the phase in (Γ complex) does a phase shift (delay) so, the Reference point changes.

• The vectorial sum at the right side equals The vectorial sum at the left side.



Continue...

POWER UNIT

$$E_i + E_r = E_t$$

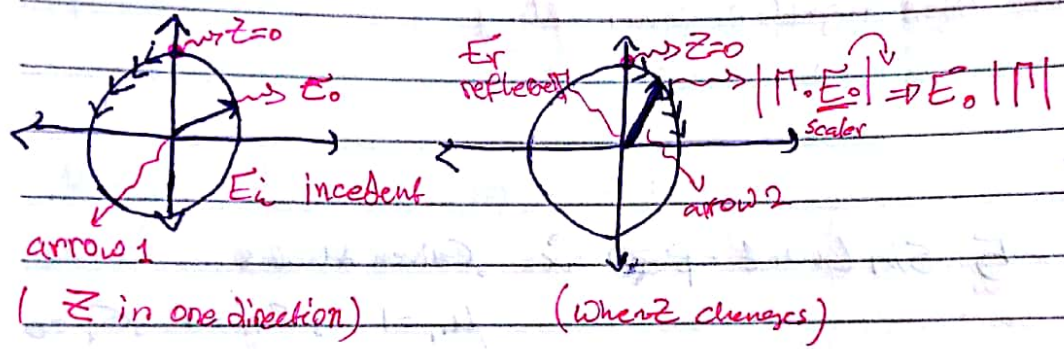
$$E_i \left(\frac{1+\Gamma}{\tau} \right) = E_t, \quad \tau \equiv \text{Transmission coefficient}$$

$$\tau E_i = E_t \quad \#$$

* First material: $E_r = \Gamma E_0 \sin(\omega t - \beta_1 z) \hat{a}_x$

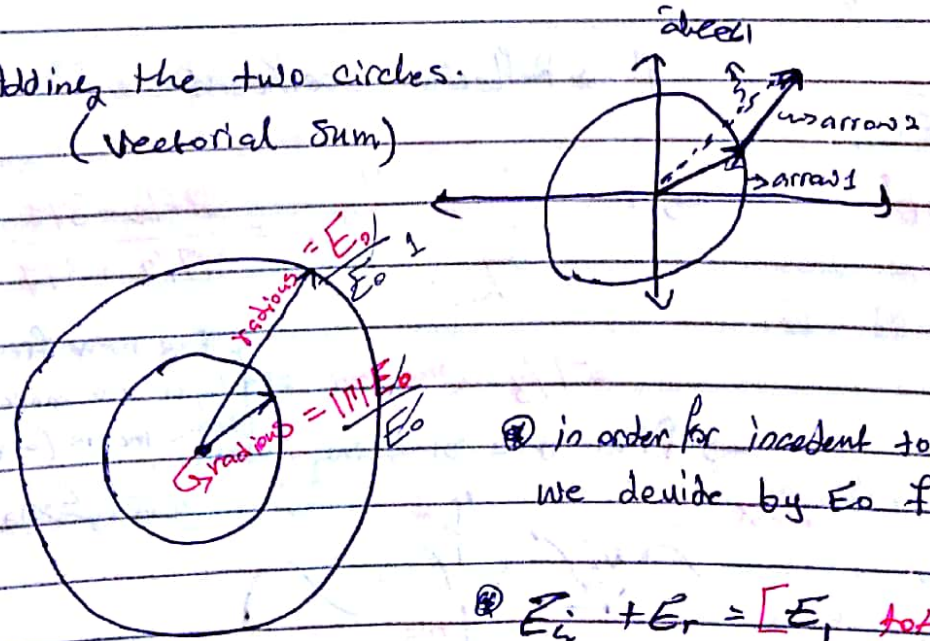
* Second material: $E_t = \tau E_0 \sin(\omega t - \beta_2 z) \hat{a}_x$

→ Note: all equations & Relations Applies on the boundary ONLY



* Rotating around the unity circle is $\lambda - 2\pi$ in the z -axis.

→ Adding the two circles: (Vectorial sum)



* in order for incident to equal 1 we divide by E_0 for simplicity

$$E_i + E_r = [E_t \text{ total in the first media}]$$

E_t is sinusoidal with double the original frequency.

time β & ω are constants and we fix time to $t = \text{zero}$ in order to remove its effect.

POWER UNIT

E_{total} :

$$E_i + E_r = E_1 \sin(\omega_0 t - \underbrace{2\beta z}_{\text{double}}) \hat{a}_x$$

Standing Waves:

called "standing" cause it will not change its position or effect. It will remain in its position at the (graph) check Figure *

Ok, then what exactly is it?
 ((two waves added together coherently))

Example

$E_i = 5 \sin(\omega_0 t - \beta z) \hat{a}_x$, Given that:

$\mu_1 = 1, \epsilon_1 = 1, \sigma_1 = 0$
 $\mu_2 = 1, \epsilon_2 = 4, \sigma_2 = 0$

Find:

$$\rightarrow \frac{u_1}{u_2} = \sqrt{\frac{\mu_1}{\epsilon_1}} = \frac{u}{u_0} = 377 \Omega$$

$$\rightarrow \frac{u_1}{u_2} = \frac{u}{2}$$

$$\rightarrow u_1 = 0 = 3 \times 10^8$$

$$\rightarrow u_2 = \frac{c}{2}$$

Γ Reflection coefficient

$$\rightarrow \Gamma = \frac{u_2 - u_1}{u_2 + u_1} = \frac{377/2 - 377}{377/2 + 377}$$

= -1/3 (negative) \rightarrow if we move from more dense media (-ve) to free space \rightarrow (50%)

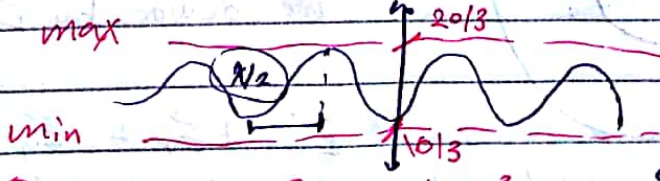
Draw the standing wave

$$\rightarrow \beta_1 = 20 \pi$$

$$\lambda_1 f = c$$

$$\lambda_1 = \frac{3 \times 10^8}{3 \text{ GHz}}$$

$$\lambda_1 = 0.1 \text{ m}$$



Maximum = $(1 + |\Gamma|) * E_0$ $\lambda/2 = 50 \text{ cm}$
 $= 4/3 * 5 = 20/3$

Minimum = $(1 - |\Gamma|) * E_0$
 $= 2/3 * 5 = 10/3$

$$\rightarrow \beta_2 = 40 \pi$$

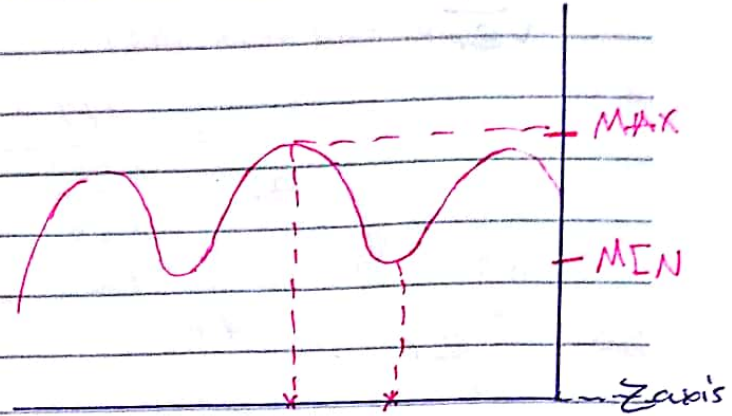
$$\lambda_2 = 5 \text{ cm}$$

POWER UNIT

Ex Given $\mu_1 = 100$, $\mu_2 = 200$, $l = 10$ cm

$$\vec{E}_i = 10 \sin(\omega t - \beta z) \vec{a}_z$$

Draw the standing wave:



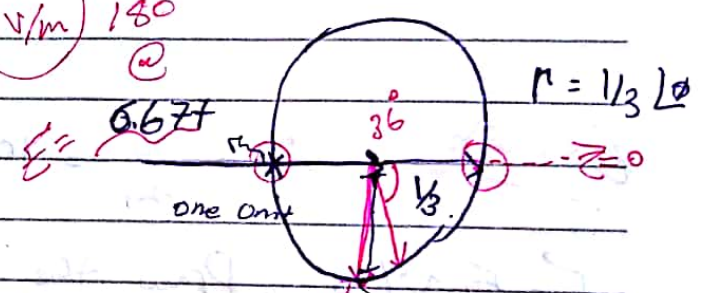
- MAX = $(1 + |\Gamma|) \times E_0$
- MIN = $(1 - |\Gamma|) \times E_0$
- $\lambda/4 =$
- $\lambda/2 =$
- $|\Gamma| = \frac{\mu_2 - \mu_1}{\mu_2 + \mu_1} = 1/3$

MAX = $(1 + \frac{1}{3}) \times 10 = 13.3$ V/m

MIN = $(1 - \frac{1}{3}) \times 10 = 6.67$ V/m

$\lambda/4 = 2.5$ cm

$\lambda/2 = 5$ cm



Find Electric Field Value @ $z = 0.5$ cm? @ 180° How on earth did we know that?

So $2.5 \rightarrow 180$
 $0.5 \rightarrow X$

$$X = \frac{180 \times 0.5}{2.5}$$

$X = 36^\circ$

Find the Value of z @ 90° (position)?
 SAME PROCESS

$2.5 \rightarrow 180$
 $X \rightarrow 90$

$$X = \frac{90 \times 2.5}{180}, \quad X = 1.25 \text{ cm}$$

Find the speed of propagation

POWER UNIT

$$u = \frac{c}{\sqrt{\epsilon_r}}$$

Find this first

we know $\frac{y}{z_1} = \frac{y}{z_0} = \frac{377 \Omega}{\sqrt{\epsilon_r}} = 100$
 from the Question

$$\text{So, } \sqrt{\epsilon_r} = \frac{377}{100}$$

$$\Rightarrow \sqrt{\epsilon_r} = 3.77 \text{ which is } \gg 1$$

$$\Rightarrow u = \frac{c}{\sqrt{\epsilon_r}} = \frac{3 \times 10^8}{3.77} = 79.57 \text{ M}$$

Example Given $\frac{y}{z_1} = 100$, $\frac{y}{z_2} = 50 + j100$, $\lambda = 12 \text{ cm}$

$E = 5 \text{ V/m}$. Draw the standing wave

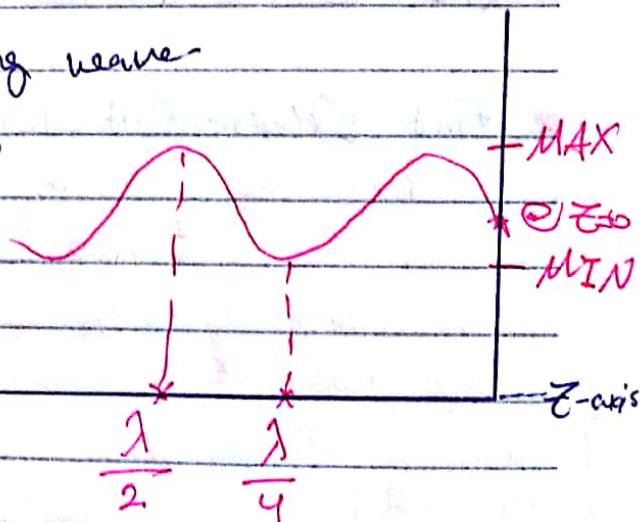
$$\Gamma = \frac{-50 + j100}{150 + j100} = 0.62 \angle 82^\circ$$

• MAX = 16.2 V/m

• MIN = 3.8 V/m

• $\lambda/4 = 3 \text{ cm}$

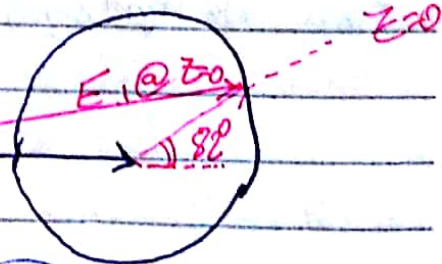
• $\lambda/2 = 6 \text{ cm}$



Find E @ $z=0$

$$= 10 * \sqrt{(1 + 0.62 \cos(82^\circ))^2 + (0.62 \cos(82^\circ))^2} \text{ [one unit]}$$

$$= 12.4 \text{ V/m}$$



Note: $\frac{1}{2} \lambda$

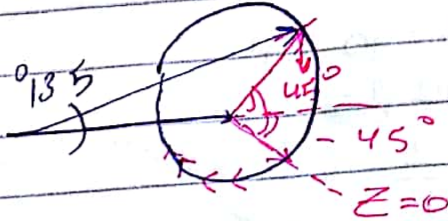
POWER UNIT

Find E @ 13.2 cm

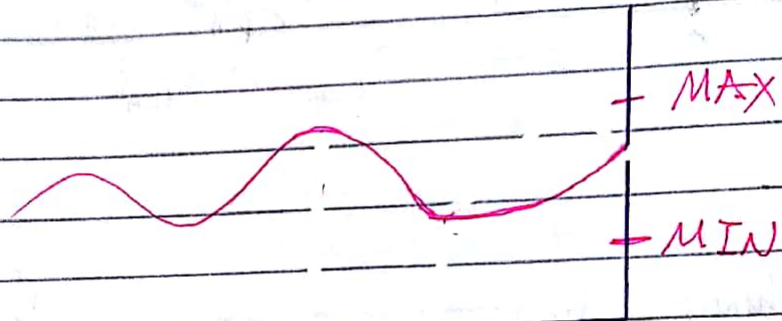
Answer

angle 75°

Ex $\Gamma = 0.5 \angle -45^\circ$
 $\lambda = 8 \text{ cm}$
 $E_0 = 4 \text{ V/m}$



- MAX = 6 V/m
- MIN = 2 V/m
- $\lambda/4 = 2 \text{ cm}$
- $\lambda/2 = 4 \text{ cm}$



⊗ Matching point is:
When the Reflection
Coefficient = Zero.

POWER UNIT

$$* E_i = 5 \sin(\omega_0 t - \beta_1 z) \vec{a}_x$$

$$* H_i = \frac{5}{\eta_1} \sin(\omega_0 t - \beta_1 z) \vec{a}_y$$

$$* E_r = \Gamma * 5 \sin(\omega_0 t \oplus \beta_1 z) \vec{a}_x$$

check?!

$$* H_r = \frac{\Gamma}{\eta_1} * 5 \sin(\omega_0 t + \beta_1 z) \vec{a}_y$$

if Γ (-ve) then $E_r \rightarrow -\vec{a}_x$, $H_r \rightarrow \vec{a}_y$

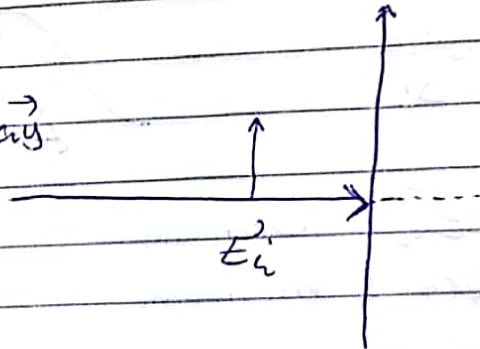
Ex → Continue the previous Example where the second material (media) :- $\Gamma = 0$, $\eta_1 = \eta_2$ (matching)

$$\tau = (1 + \Gamma)$$

So $\tau = 1$ & $E_i = E_t \Rightarrow$ wave will continue straight

$$E_t = 5 \sin(\omega_0 t - \beta_2 z) \vec{a}_x$$

$$H_t = \frac{5}{\eta_2} \sin(\omega_0 t - \beta_2 z) \vec{a}_y$$



if the second material

have $\mu = 2\mu_0$, $\epsilon = 2\epsilon_0$

then...

$\eta = \eta_0$ & β is doubled

but the speed $u = \frac{c}{2}$

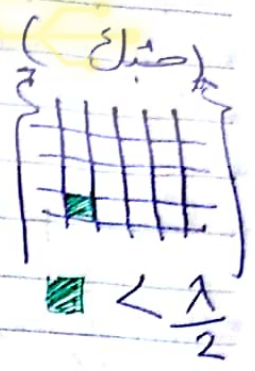
Which means that you can change the speed of propagation and also change the material but still have matching.



if we had a conductor $\eta_2 = 0$, $\Gamma = -1$

POWER UNIT

If the second material is a conducting grid and the openings are smaller than $\frac{\lambda}{2}$, then we say that the grid is a good conductor (short circuit)



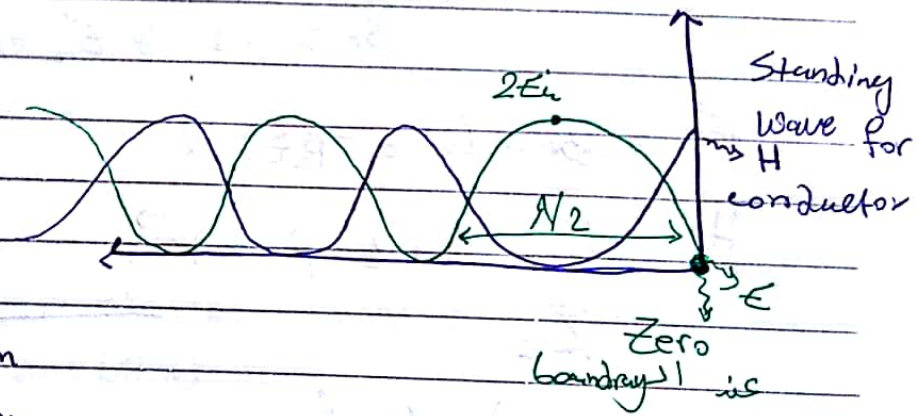
$\Gamma = 1$

$\Gamma = -1$

this means that we are coming out of the conductor toward the non-conducting material. non-realistic as to conductor propagation.

this means that we are entering the conductor. all "live" wires Electric field will be zero. Real case ↑

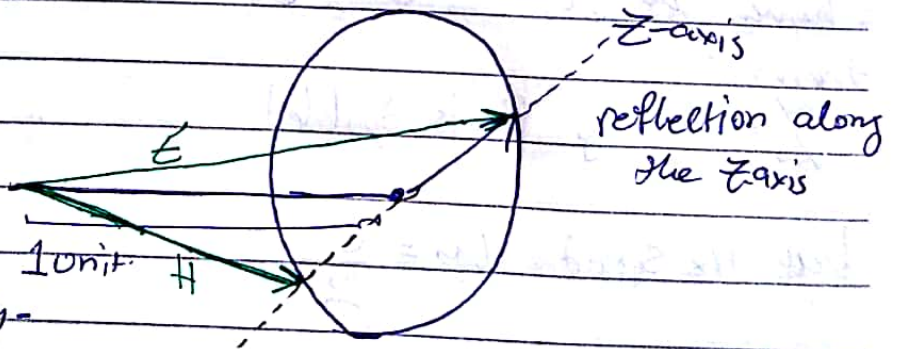
$H = \omega \epsilon_0 E$



$2E_i \equiv \text{maximum}$
 $Zero \equiv \text{minimum}$

$E = \frac{1}{n} H$

When E is max H is min & vice versa.



Homework \Rightarrow ReDo all Examples but this time include H Drawing & calculate its value

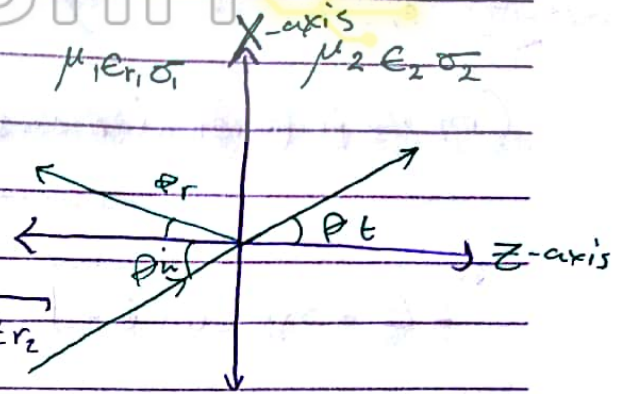
⊗ Oblique Incident

$$\theta_r = \theta_i$$

From Snell's law

$$n_1 \sin(\theta_i) = n_2 \sin(\theta_t)$$

$$\text{where } n_1 = \sqrt{\mu_1 \epsilon_1}, \quad n_2 = \sqrt{\mu_2 \epsilon_2}$$



Two General cases are

⊗ (perpendicular case) ⊗

$$1) \quad \Gamma_{\perp} = \frac{\frac{\mu_2}{\epsilon_2} \cos \theta_i - \frac{\mu_2}{\epsilon_2} \cos \theta_t}{\frac{\mu_2}{\epsilon_2} \cos \theta_i + \frac{\mu_2}{\epsilon_2} \cos \theta_t}$$

perpendicular on the plane that has the trace for two materials

$$2) \quad \tau_{\perp} = \frac{2 \frac{\mu_2}{\epsilon_2} \cos \theta_i}{\frac{\mu_2}{\epsilon_2} \cos \theta_i + \frac{\mu_2}{\epsilon_2} \cos \theta_t}$$

⊗ (parallel case) ⊗

Incident Electric Field

$$1) \quad \Gamma_{\parallel} = \frac{\frac{\mu_2}{\epsilon_2} \cos \theta_t - \frac{\mu_1}{\epsilon_1} \cos \theta_i}{\frac{\mu_2}{\epsilon_2} \cos \theta_t + \frac{\mu_1}{\epsilon_1} \cos \theta_i}$$

$$2) \quad \tau_{\parallel} = \frac{2 \frac{\mu_2}{\epsilon_2} \cos \theta_i}{\frac{\mu_2}{\epsilon_2} \cos \theta_t + \frac{\mu_2}{\epsilon_2} \cos \theta_i}$$

⊗ The point is to find E_r (reflected)

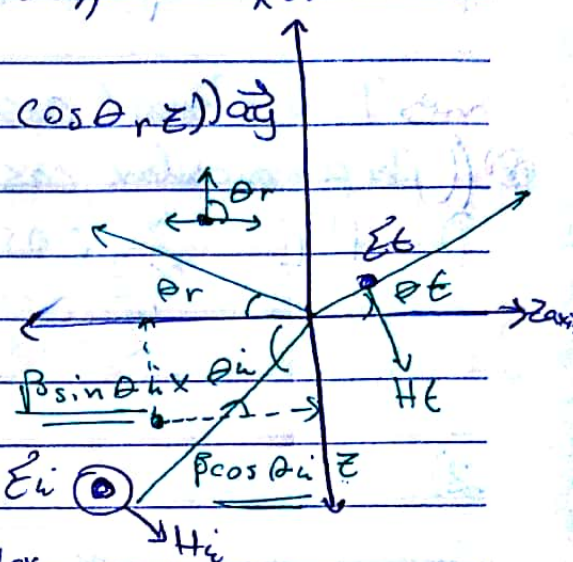
$$(\Gamma_r \epsilon_i)$$

Example $\mu_{r1} = \mu_{r2} = 1 \parallel \epsilon_{r1} = 1, \epsilon_{r2} = 4 \parallel \sigma = 0$

(on the perpendicular case) $E_i = E_0 \sin(\omega t - \beta r) \vec{a}_y$
 $\underline{\underline{\beta_r}} = (\sin \theta_i \vec{a}_x + \cos \theta_i \vec{a}_z)$

$$= E_0 \sin(\omega t - \beta_1 (\sin \theta_i x + \cos \theta_i z))$$

$$E_r = \Gamma_{\perp} E_0 \sin(\omega t - \beta_1 (\sin \theta_r x - \cos \theta_r z)) \vec{a}_y$$



$$\rightarrow n_1 \sin \theta_i = n_2 \sin \theta_t$$

$$\rightarrow n_1 = \sqrt{\mu_{r1} \epsilon_{r1}}, n_2 = \sqrt{\mu_{r2} \epsilon_{r2}}$$

$$\rightarrow \sin \theta_i = 2 \sin \theta_t$$

⊗ The second media is more dense so that θ_t (with z-axis) is smaller

$$\epsilon_{r2} > \epsilon_{r1}$$

Using $\sin \theta_t = \frac{n_1}{n_2} \sin \theta_i = 1 \Rightarrow \sin^{-1} \left(\frac{n_2}{n_1} \right)$

$\frac{n_2}{n_1} \geq 1$, if $\theta_i = 30$ in this Example.

$$\begin{aligned} & \hookrightarrow \sin^{-1} \left(\sin(30) \cdot 2 \right) \\ & = \sin^{-1} \left(\frac{1}{2} \cdot 2 \right) \Rightarrow 14.5 \end{aligned}$$

$$E_t = E_0 T_{\perp} \sin(\omega t - \beta_2 (\sin \theta_t x + \cos \theta_t z)) \vec{a}_y$$

$$H_i = \frac{E_0}{Z_1} \sin(\omega t - \beta_1 (\sin \theta_i x + \cos \theta_i z)) \cdot (\text{unit vector})$$

Unit vector $\vec{a}_i = (-\cos \theta_i \vec{a}_x + \sin \theta_i \vec{a}_z)$

$$H_r = \frac{E_0 \Gamma_{\perp}}{Z_1} \sin(\omega t - \beta_1 (\sin \theta_r x + \cos \theta_r z)) \cdot (\text{unit vector})$$

(Unit vector $\vec{a}_r = (\cos \theta_r \vec{a}_x + \sin \theta_r \vec{a}_z)$)

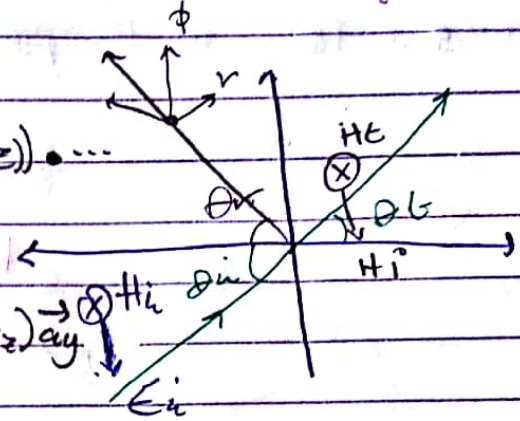
$$H_L = \frac{E_0 \tau_{\perp}}{\mu_0} \sin(\omega t - \beta_2 (\sin \theta_L x + \cos \theta_L z)) \cdot (\text{unit vector})$$

unit vector $\Rightarrow (-\cos \theta_L \hat{a}_x + \sin \theta_L \hat{a}_z)$

Homework Redo the Example in the parallel case

Note \Rightarrow μ_0 μ_0

$$\rightarrow E_i = E_0 \sin(\omega t - \beta_1 (\sin \theta_i x + \cos \theta_i z)) \cdot (\cos \theta_i \hat{a}_x - \sin \theta_i \hat{a}_z)$$



$$H_i = \frac{E_0}{\mu_0} \sin(\omega t - \beta_1 (\sin \theta_i x + \cos \theta_i z)) \hat{a}_y$$

$$\rightarrow E_L = E_0 \tau_{\parallel} \sin(\omega t - \beta_2 (\sin \theta_L x + \cos \theta_L z)) \cdot (\cos \theta_L \hat{a}_x - \sin \theta_L \hat{a}_z)$$

$$H_L = \frac{E_0}{\mu_0} \tau_{\parallel} \sin(\omega t - \beta_2 (\sin \theta_L x + \cos \theta_L z)) \hat{a}_y$$

$$\rightarrow E_r = E_0 \Gamma_{\parallel} \sin(\omega t - \beta_1 (\sin \theta_r x - \cos \theta_r z)) \cdot (-\cos \theta_r \hat{a}_x - \sin \theta_r \hat{a}_z)$$

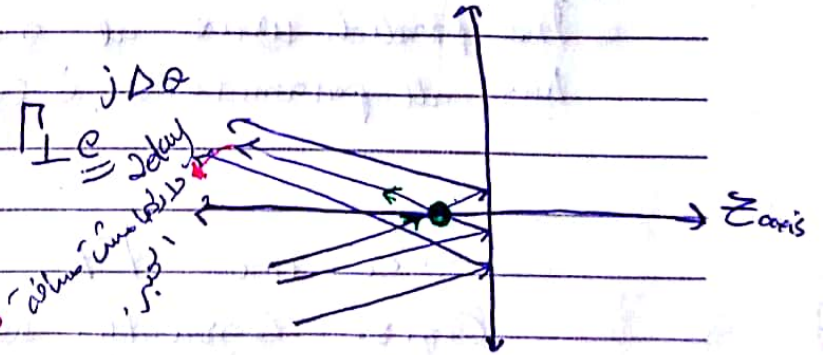
$$H_r = \frac{E_0}{\mu_0} \Gamma_{\parallel} \sin(\omega t - \beta_1 (\sin \theta_r x - \cos \theta_r z)) \hat{a}_y$$

Standing Wave

$$\theta_0 = \theta_{\perp} + \Delta \theta = \Delta \theta \perp e^{j\Delta \theta}$$

$$E_r \hat{a}_y = E_0 \Gamma_{\perp} e^{j\Delta \theta}$$

so $\omega t - \beta z$ is the same



Wave front of electric field

are coherent in phase of fields

incident wave is reflected wave. They are at the same frequency and the same plane/space.

POWER UNIT

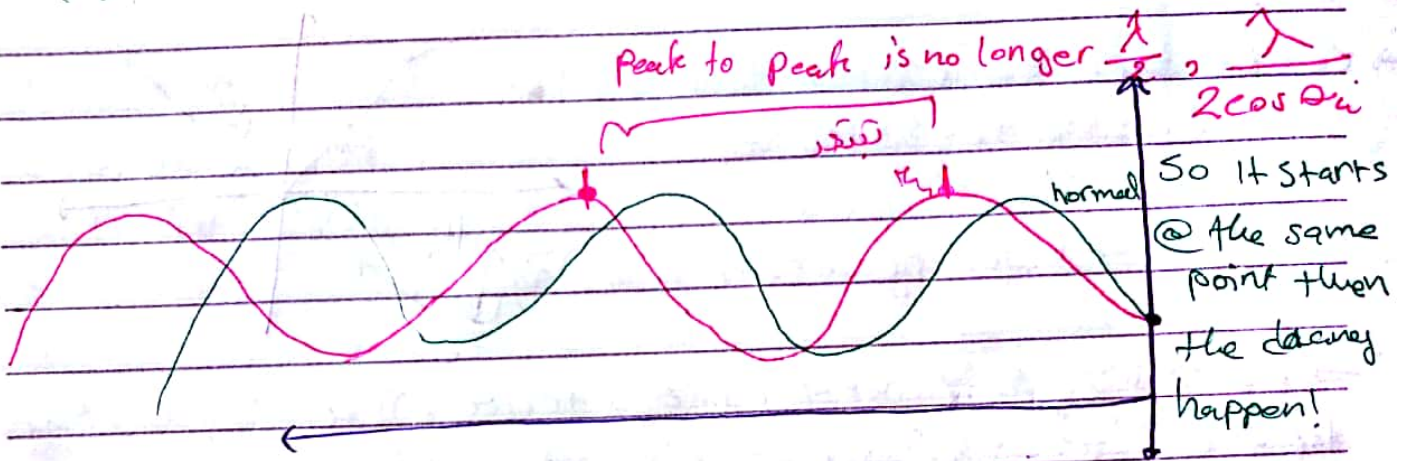
جزء اولاً Realativity as just

$$\Delta \theta = 2\beta Z \cos \theta_i$$

Zero $\theta_0 = \theta_{\pi}$ in the normal incidence

$$\theta_i = \theta_{\pi} + \Delta \theta$$

$\theta_0 = \theta_0$ in the parallel case [No Reflection]



$Z=0$, $\theta_0 = \theta_{\pi}$, الكمية الدوران تتغير عند Z

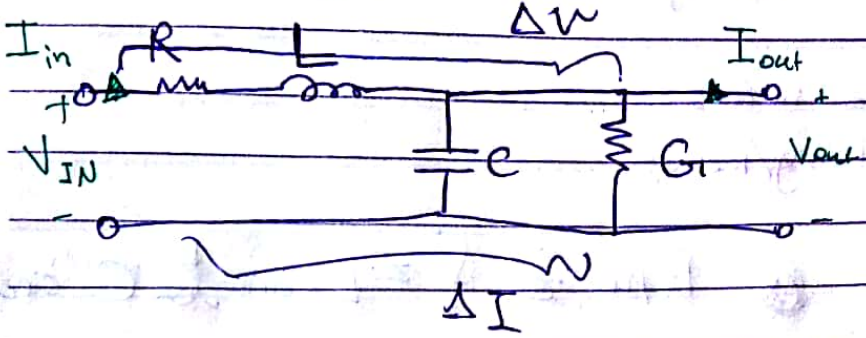
Source $Z=0$ wave and $Z=0$ wave
 plane and $Z=0$ minimum

The parallel wave at a constant Z has all parameters decreasing!

Lecture 17 is not written
 cause basically we didn't learn anything
 important but! if you're unpleased with
 that fact, the record for the lecture
 will be available. Go check it on your
 own!

Transmission Line

one Unit equivalent circuit model for transmission line.



$$V_{out} = \Delta V$$

$$I_{out} = \Delta I$$

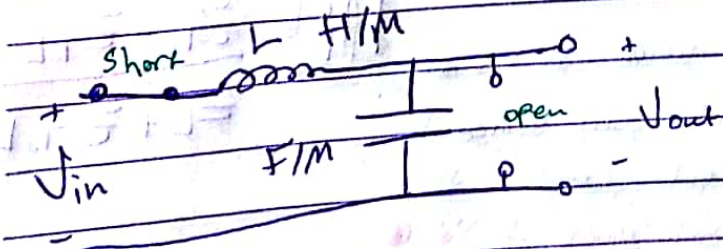
$$\nabla^2 V - \gamma V = 0$$

$$V = V_0 \sin(\omega t - \beta x)$$

No direction

دیس کی طرف سے

Lossless Transmission Line



$$C = \infty \quad R = 0$$

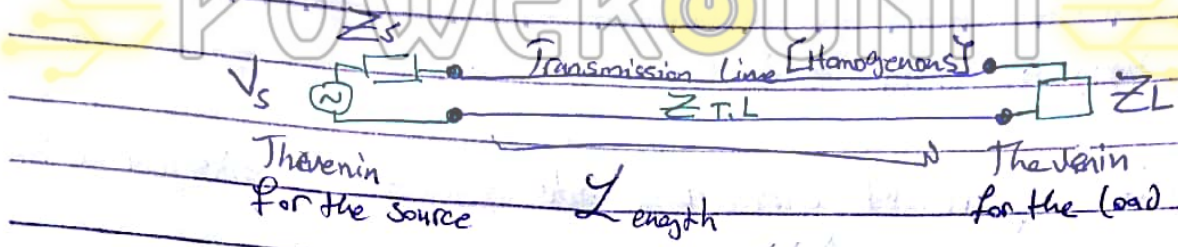
- Characteristic Impedance
- Speed of propagation in TL

$$Z_{TL} = \sqrt{\frac{L}{C}}$$

$$u = \frac{1}{\sqrt{L \cdot C}}$$

$$f \cdot \lambda = u$$

POWER UNIT

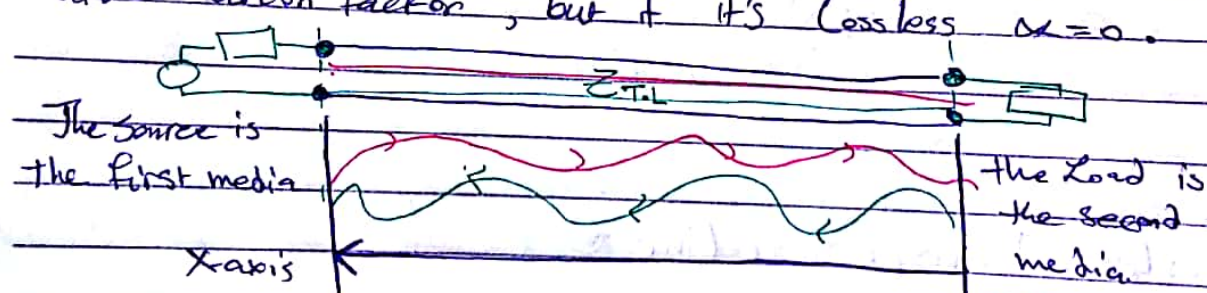


Small $L \equiv$ meters
 $L \equiv$ Henry/meter

Beware!!
 $Z \neq L$
 (استبدلها في الـ Z بالـ L)

- Both Z_s & Z_L are LC circuits
- $Z_{T.L}$ is Real if it's lossless

attenuation factor, but if it's lossless $\alpha = 0$.



$$\Gamma_s = \frac{Z_s - Z_{T.L}}{Z_s + Z_{T.L}}$$

Lead of wave at source

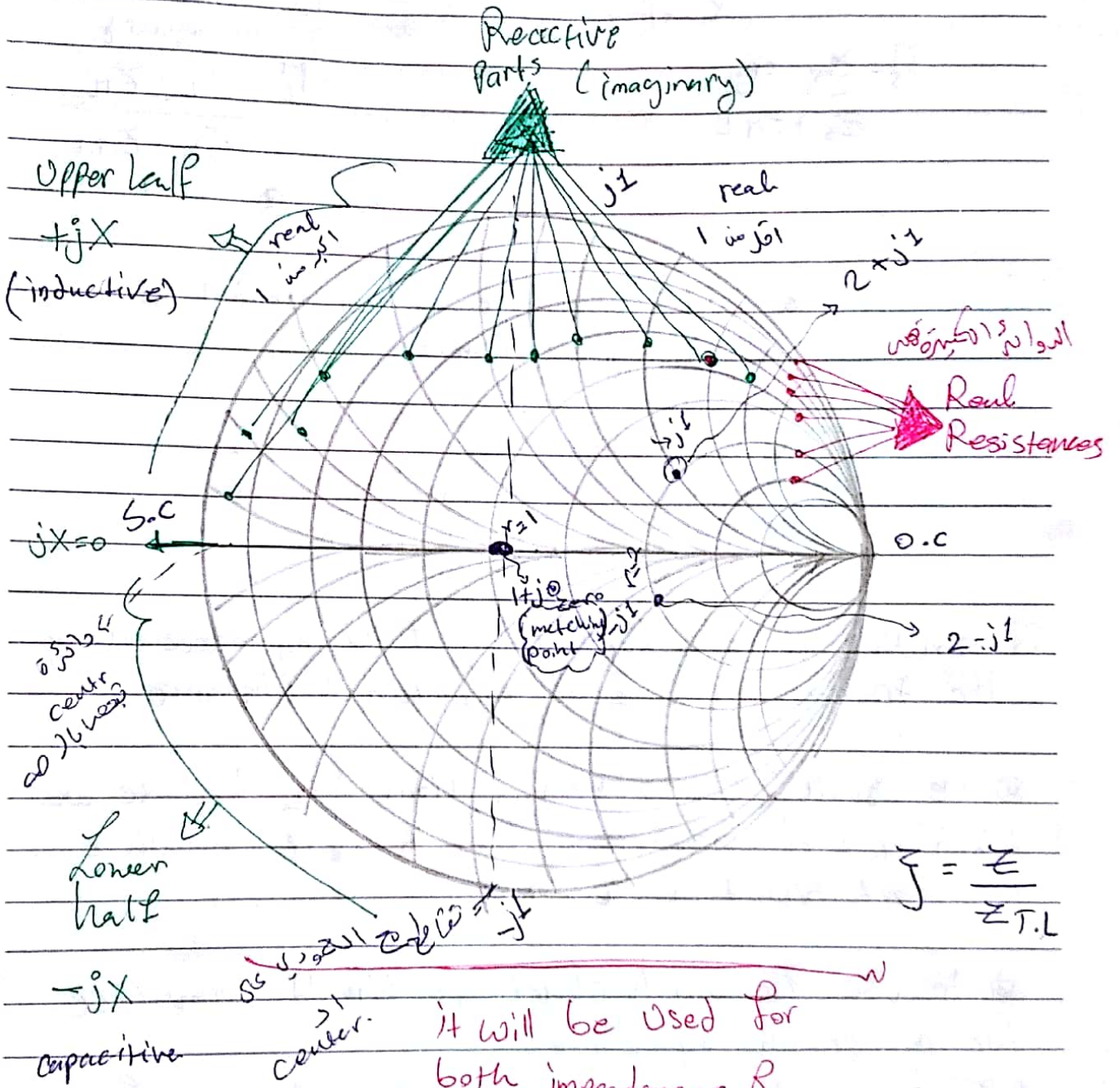
$$\Gamma_L = \frac{Z_L - Z_{T.L}}{Z_L + Z_{T.L}}$$

Lead of wave at load

- Maximum power transfer is when the wave (as a whole) is sent from the source to the load without reflecting back!
- We don't like waves/signals that reflect
- To avoid reflection, $\Gamma = 0$ and this happens when $Z_L = Z_{T.L}$, called Matching.

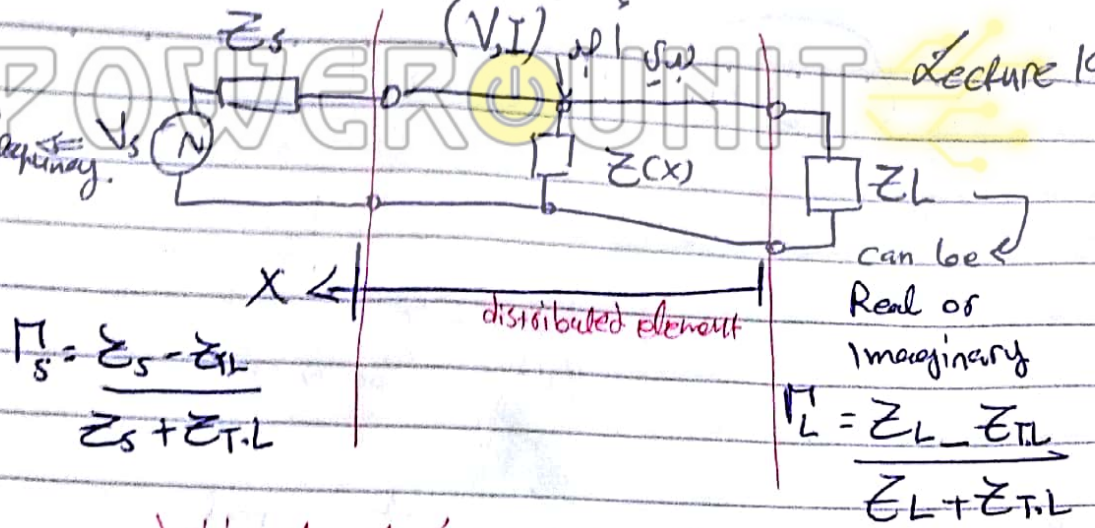
Smith Chart

الدائرة الخيالية المثلثة



الدائرة الخيالية المثلثة
center
center
it will be used for both impedances & admittances

High Frequency V_s



$$\Gamma_s = \frac{Z_s - Z_{T.L}}{Z_s + Z_{T.L}}$$

$$\Gamma_L = \frac{Z_L - Z_{T.L}}{Z_L + Z_{T.L}}$$

distributed element - لا يمكن ان نتعامل معها كعنصر موزع

وتعتبر كعنصر موزع

⊗ 0.1λ and above we will treat it as a distributed element.

⊗ distribution of voltage & current is cyclic so to normalize things
 توزيع موج ال T.L ال 0.1λ ←

⊗ Smith chart is a Representation of the impedance along the transmission line as a function of distance

⊗ we usually starts analysis from Z_L because we know alot of information about it, but anywhere we start should be the same.

⊗ to use the Smith chart we should normalize for a certain impedance.

→ we normalize by using a normalization factor
 → Length of the T.L $Z_{T.L}$ (base impedance) Z_0 should also be normalized to λ

⊗ capacitors and inductors will cause delays if the load in the second media isn't matched.

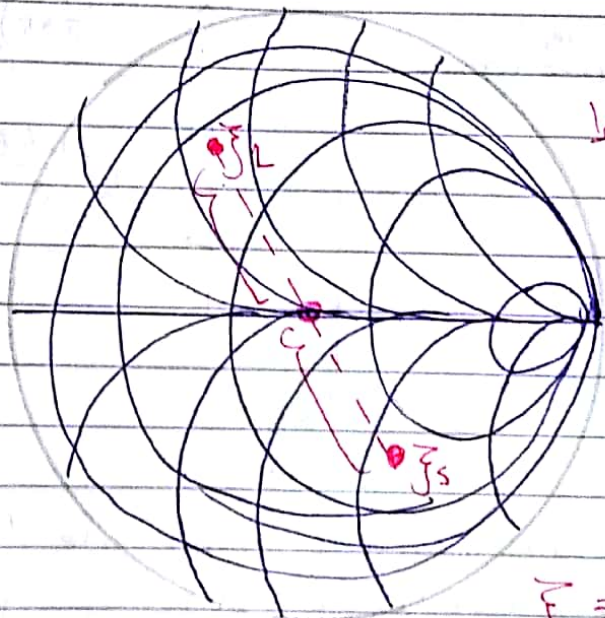
⊙ After delays & Reaching the Second media we will have Reflection.

⊙ Transient State's I can only see the effect in a specific time instant.

⊙ Steady States I see the effect every where along the Transmission Line @ any time

⇒ THE SMITH CHART IS ONLY FOR STEADY STATE ANALYSIS.

$$\Gamma = \frac{Z - Z_{T.L}}{Z + Z_{T.L}}$$

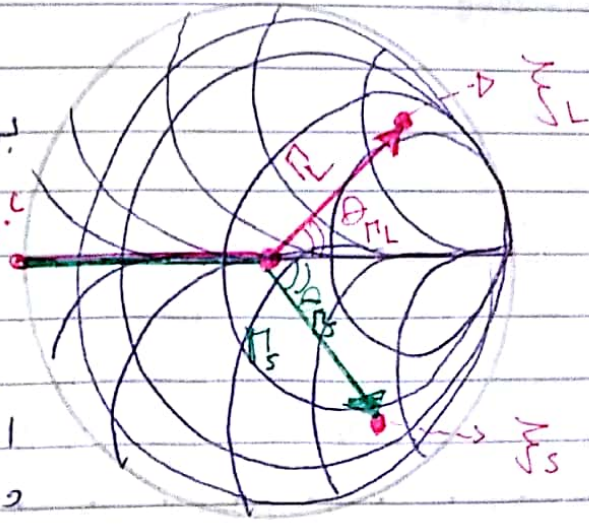


Ⓣ اذا كان Γ_L في المحور
 والـ Γ_S في
 Center الـ Γ_S
 ووجه الـ Γ_S
 وكذا الـ Γ_L
 بالـ Γ_S
 الـ Γ_S

$$\Gamma_S = \frac{1}{\Gamma_L}$$

Ⓣ If I wanted to Draw $|\Gamma_L| \angle \theta_{\Gamma_L}$ as a complex #

- ① رسم دائرة الـ S.C الـ Center الـ S.C
- ② رسم الـ Vector الـ $|\Gamma_L| \angle \theta_{\Gamma_L}$ بالـ Γ_S
- ③ الزاوية بين الـ Vectors الـ Γ_L و Γ_S الـ θ_{Γ_L}

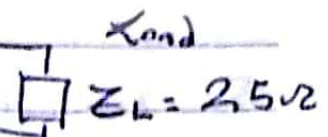
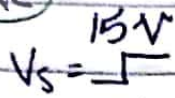


$$|\Gamma_S| \angle \theta_{\Gamma_S}$$

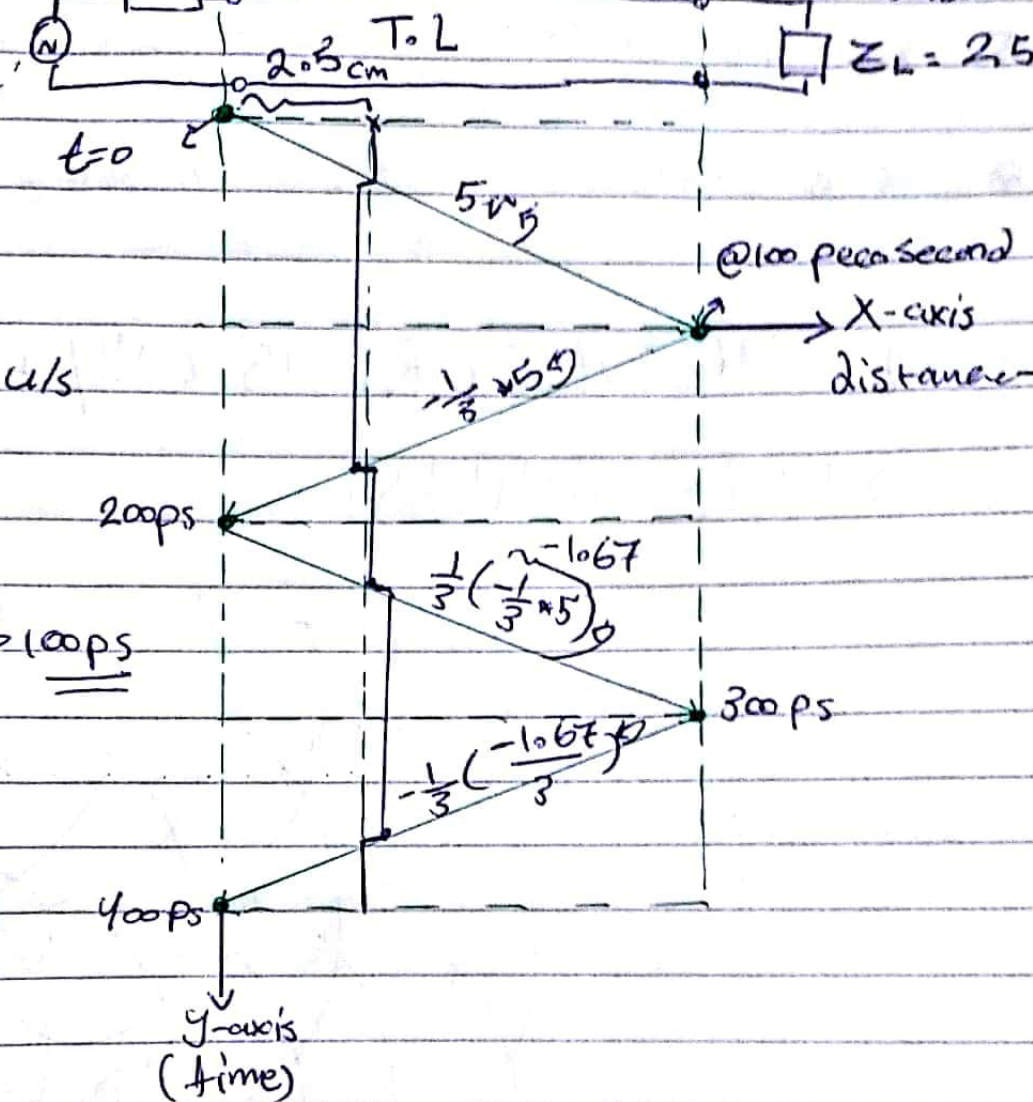
Transient Analysis

Lecture 30

Example



not a sinusoidal voltage
(Unit Step)



$$u = \frac{c}{3} = 1 \times 10^8 \text{ u/s}$$

$$u = \frac{\text{distance}}{\text{Time}}$$

$$\Rightarrow \text{Time} = 1 \times 10^{-10} \Rightarrow \underline{\underline{100 \text{ ps}}}$$

$$\Gamma_{\text{load}} = -\frac{1}{3}$$

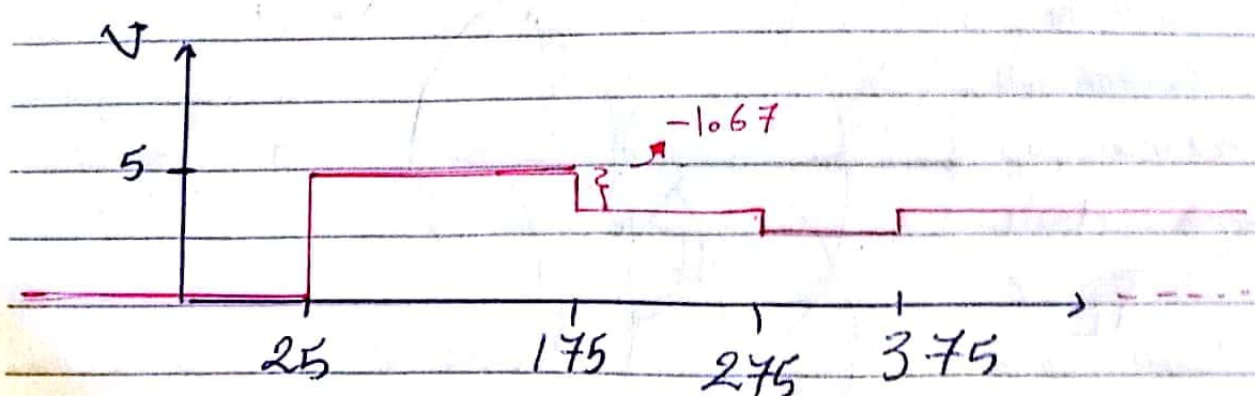
$$\Gamma_{\text{source}} = \frac{1}{3}$$

Every time we reach a media we must multiply by its Reflection Coefficient (Γ)

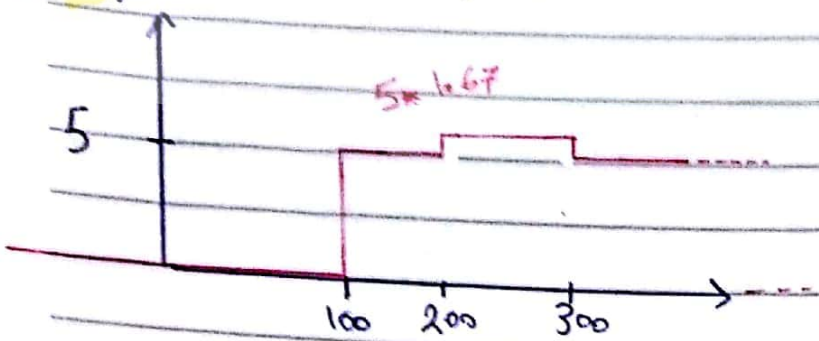
Γ Wave 100%
! Reflection will be 100% at 5

Draw the Voltage

@ 2.5 cm



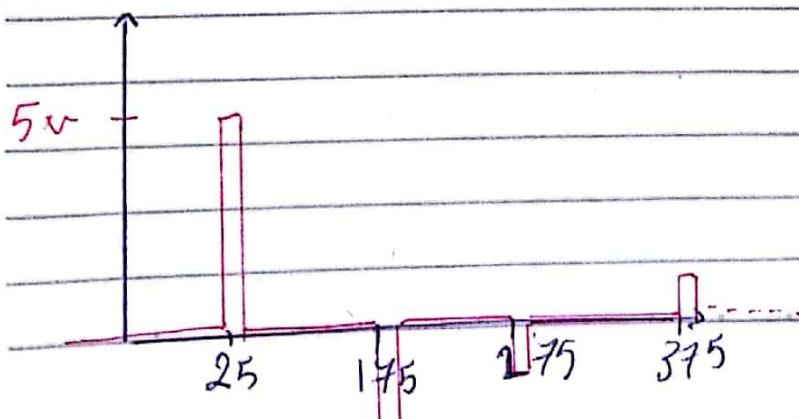
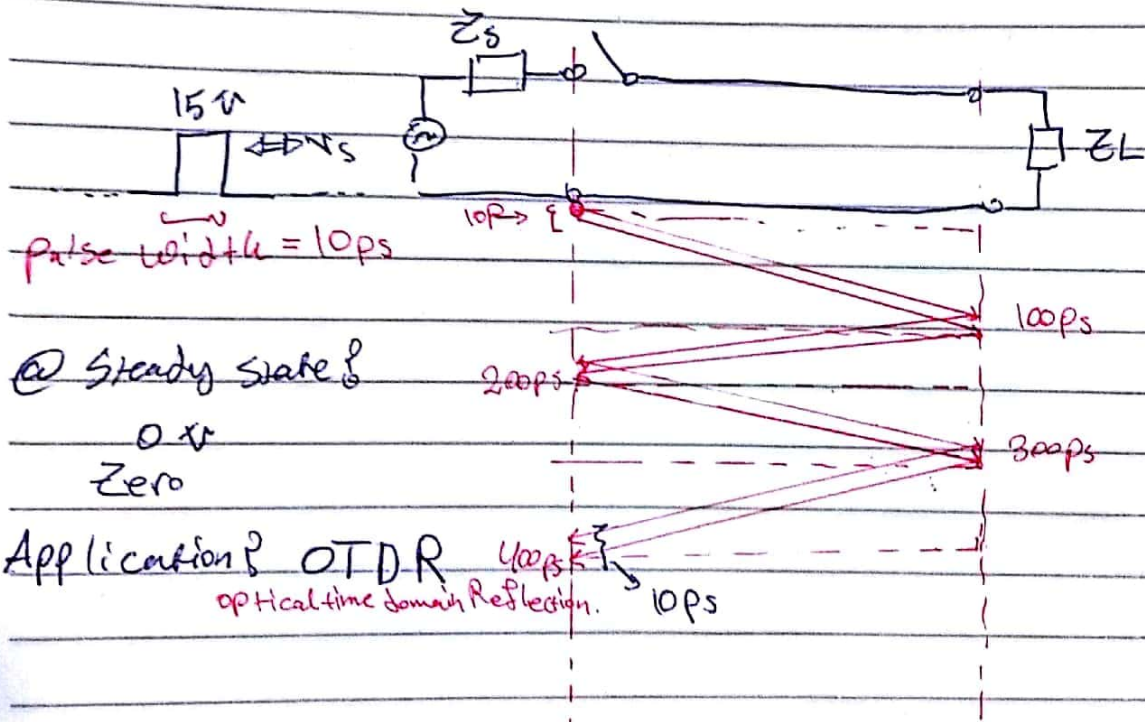
Draw the Voltage @ the Load



@ steady state [OR] after a long time the voltage will be @ the Load

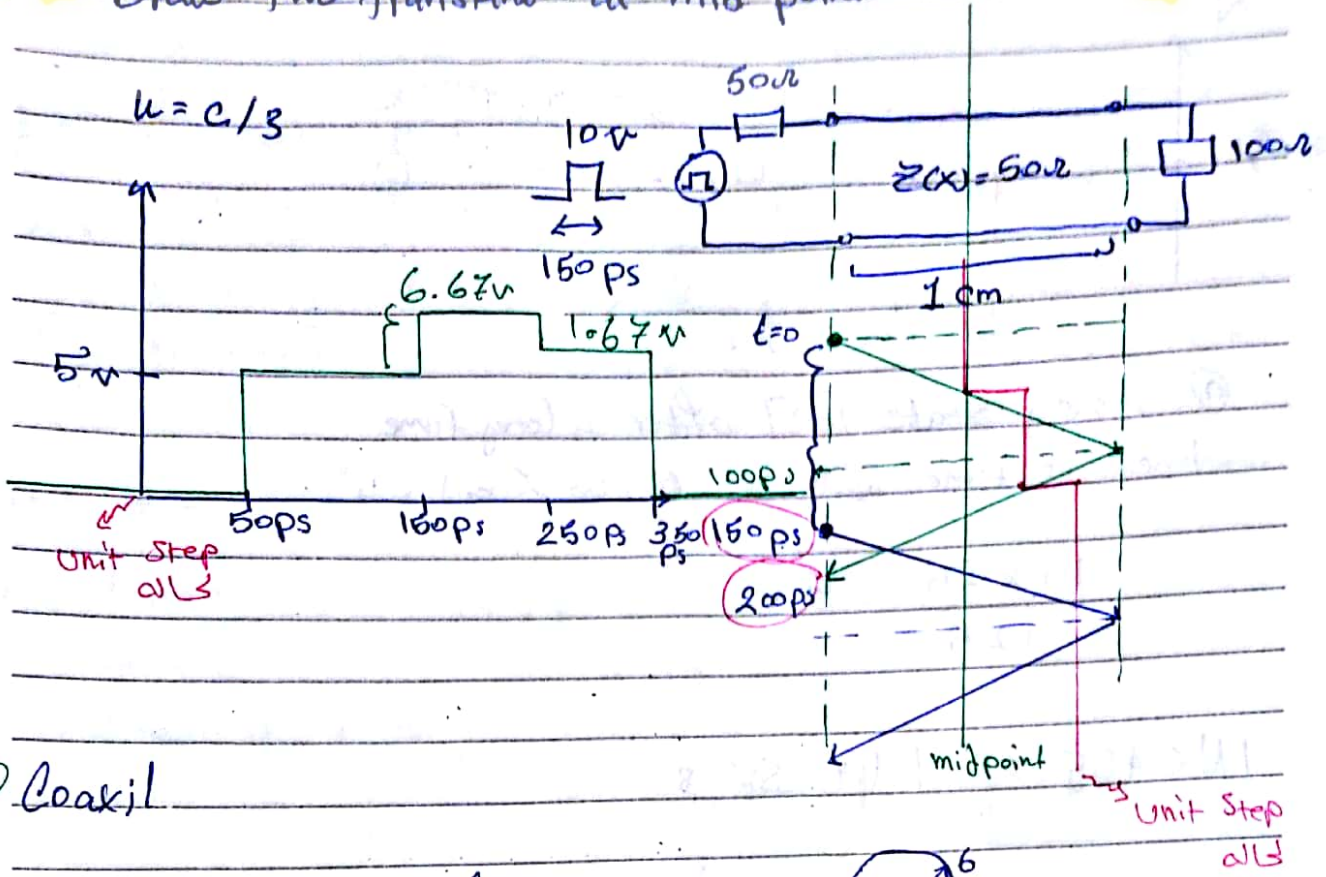
$$\frac{15 \times 25}{125} = 3 \text{ v}$$

IN CASE OF PULSE



Quiz

Draw the transient at mid point



Coaxial

$$Z_{T.L} = \frac{138}{\sqrt{\epsilon_r}} \cdot \ln\left(\frac{b}{a}\right), \text{ where } \begin{matrix} b \\ \circlearrowleft \\ a \end{matrix}$$

$$Z_{T.L} = \frac{276}{\sqrt{\epsilon_r}} \cdot \ln\left(\frac{b}{a}\right), \text{ where } \begin{matrix} b \\ \circlearrowleft \\ a \end{matrix}$$

Micro Strip

$$Z_{T.L} = \frac{377}{\sqrt{\epsilon_r} \cdot \left(\frac{\text{Width}}{\text{Height}} + 2\right)}, \text{ where } \begin{matrix} \text{Width} \\ \text{Height} \end{matrix}$$

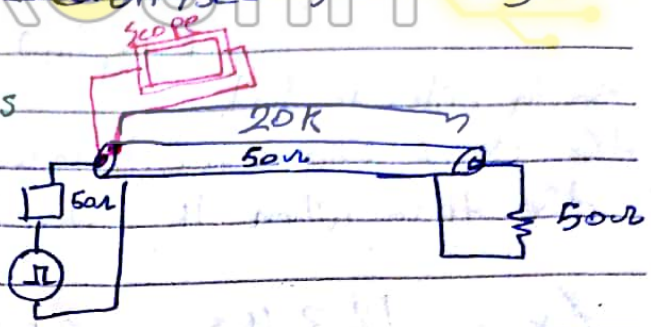
Total Delay = $\frac{\text{Length}}{\text{Speed}}$

$v = \frac{1}{\sqrt{L \cdot C}}$
Speed of propagation

Ex assume DataRate : 1 Gbit/sec, $\mu = 0.13$

Delay: $= \frac{20 \times 10^3}{1 \times 10^8} = 200 \mu s$

Number of Stored bits = $\frac{200 \times 10^{-6}}{1 \times 10^{-9}} = 200k \text{ bit}$



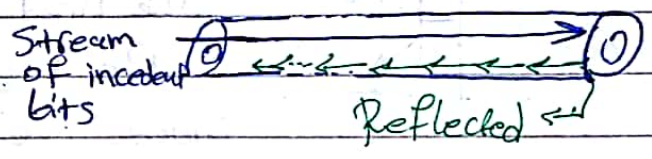
because \leftarrow

$\frac{1 \text{ bit}}{T_0} = 1 \text{ nanoSec}$

The previous type of channel is called \Rightarrow Storage device \ominus channel with memory.

Solution for the following problems caused by a channel with memory \Rightarrow ((Matching))

First problem



((Long Distances))

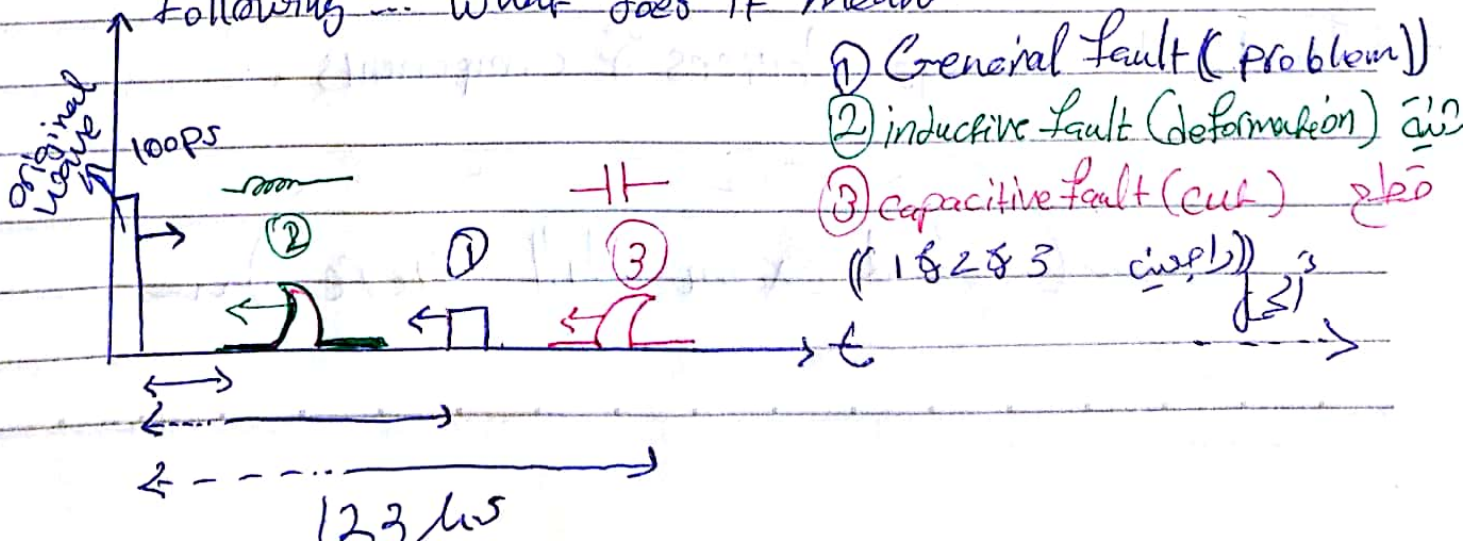
Inter symbol interference

Second problem

((Short Distances))

Fading effect

If we added a **Scope** and we found the following ... what does it mean



- ① Coercional Fault (problem)
- ② inductive fault (deformation) $\bar{a}i^2$
- ③ capacitive fault (cut) $\bar{a}i^0$

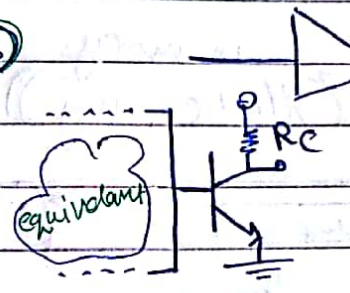
POWER UNIT

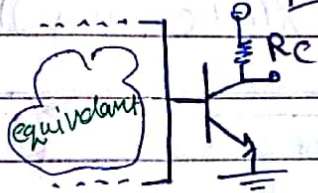
⊗ Now (123 μ s) \rightarrow Round trip

So in order to find the exact location (distance) for the fault we need the time when it started to reflect.

$$t_x = \frac{123 \mu\text{s}}{2} = 61.5 \mu\text{s}$$

So $L_x = t_x \times u^{\text{speed}}$
 $= 61.5 \mu\text{s} \times \frac{c}{3} = 6.150 \text{ m}$

⊗  Amplifier what is the o/p impedance?



we need the o/p impedance ($R_c + \text{equivalent}$) to equal 50Ω since the load is 50Ω to match up with it and avoid reflection.

⊗ if the voltage reflection coefficient is 0.1 then the reflected POWER is 100 W

$$\Rightarrow \text{Power} = \frac{(V)^2}{R}$$

\Rightarrow This will damage devices & components.

Mid exam till here